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STRONG BLAST WAVE COMPUTER PROGRAMS

Aivars Celmiņš

September 1980



**US ARMY ARMAMENT RESEARCH AND DEVELOPMENT COMMAND**  
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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) jmk This report describes a computer program package for the computation of the flow field within a strong blast bubble. The programs are based on Sedov-Laporte-Chang formulas and compute any of the following: shock front location and corresponding flow values, flow profiles at fixed times, flow histories at fixed distances, particle trajectories, and Mach-lines. The input and the output are in terms of dimensional quantities expressed in SI base units. The flow field can be calculated for a ratio of specific heats between one and seven, and for spherical, cylindrical, or planar symmetry.		

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## 1. INTRODUCTION

Strong blast waves are produced in gaseous media by a sudden deposition of large amounts of energy in a relatively small region. Such events may be approximated mathematically by a point source wave. If one assumes in addition an ideal gas and an initial pressure negligible compared to the shock pressure, then the flow field can be described in closed form as a self-similar solution of the flow equations<sup>1,2,3</sup>. A derivation of such closed form solution, including spherical, cylindrical and planar blast waves, is given, e.g., by Sedov<sup>4</sup> and Rouse<sup>5</sup>. The solutions have been supplemented by Laporte and Chang<sup>6,7</sup> who derived closed form expressions also for the particle paths and for the Mach lines.

---

<sup>1</sup>H.A. Bethe, K. Fuchs, J. von Neumann, R. Peierls, and W.G. Penny, U.S. Atomic Energy Commission Report AECD-2860 (1944).

<sup>2</sup>J.L. Taylor, "An Exact Solution of the Spherical Blast Wave Problem", *Phil. Mag.*, Vol 46 (1955).

<sup>3</sup>R. Latter, "Similarity Solution for a Spherical Shock Wave", *Journal of Applied Physics*, Vol 26, No. 8 (1955).

<sup>4</sup>L.I. Sedov, "Similarity and Dimensional Methods in Mechanics", Academic Press, New York (1959).

<sup>5</sup>C.A. Rouse, "Theoretical Analysis of the Hydrodynamic Flow in Exploding Wire Phenomena, in "Exploding Wires", W.G. Chace and K.M. Howard, eds., Plenum Press, New York, pp. 227-263 (1959).

<sup>6</sup>O. Laporte and T.S. Chang, "Exact Expressions for Curved Characteristics Behind Strong Blast Waves", U.S. Army Ballistic Research Laboratory Contractor Report, BRL-CR-30 (January 1971). (AD#722777)

<sup>7</sup>O. Laporte and T.S. Chang, "Curved Characteristics Behind Blast Waves", *Physics of Fluids*, Vol 15, pp. 502-504 (1972).

The closed form solutions are valuable approximations of physical phenomena, and they have been used as test cases for numerical solvers of flow governing equations<sup>8,9</sup>. The evaluation of the rather cumbersome formulas is, however, not trivial because the solution is expressed in parameter form. Therefore, also several tables of strong blast wave solutions have been published<sup>4,8,10,11</sup>. The use of such tables is, of course, limited to the particular set of parameters chosen by the authors of the tables. Also, because tables are expressed in form of dimensionless variables, they must be transformed before application to suit each particular case. Not all authors of the published tables provide sufficient and clear instructions as to how the transformations should be accomplished, nor are the tables of various authors standardized. Finally, the particle path and Mach line solutions of Laporte and Chang have not been tabulated at all.

In order to make the closed form solutions of strong blast waves easily available, we have developed a package of computer programs that produces the solution for any dimension and combination of parameters, within limits. The present report is a description of the programs and of their use.

In Section 2 of this report we shall summarize the formulas needed for flow computation. We shall not give a derivation of the formulas which can be obtained, e.g., from References 4 and 6, and which is based on the following assumptions:

1. the flowing medium is an ideal gas;
2. the initial pressure in the ambient gas can be neglected;

---

<sup>8</sup>P.C. Chou and R.R. Karpp, "Solution of Blast Waves by the Method of Characteristics", Drexel Institute of Technology Report, DIT Report No. 125-7, (1965).

<sup>9</sup>J.A. Schmitt, "A Finite Element Method and Corresponding Pilot Computer Code for Hyperbolic Systems of Equations in Two Spatial Dimensions and Time Applied to Unsteady Gas Flows", U.S. Army Ballistic Research Laboratory Report, BRL-R-2017 (September 1977). (AD#A045703)

<sup>10</sup>N. Gerber and J.M. Bartos, "Tables of Cylindrical Blast Functions for  $\gamma = 5/3$  and  $\gamma = 7/5$ ", U.S. Army Ballistic Research Laboratory Report, BRL-R-1376 (October 1961). (AD#663821)

<sup>11</sup>J.W. Goresch and R.G. Dunn, "Tables of Blast Wave Parameters, I Spherical Explosions", Aerospace Research Laboratory Report, ARL-69-0011 (January 1969).



3. the energy is released instantaneously at a point, in a line or in a plane.

For applications of the results to real life situations, the assumptions mean that the formulas may describe accurately events that are neither too far from the explosion (assumption 2), nor too close to it (assumption 3).

In Section 3 we shall describe the use of the computer programs for strong blast. Some examples of calculations will be presented in Section 4.

In Section 5 we derive quantitative conditions under which the strong blast solutions can provide approximate information about pressures generated by real life explosions in air.

## 2. THEORETICAL BACKGROUND

### 2.1. Shock Formulas

Let  $x_s$  be the location of the shock, i.e., the distance of the shock from the explosion, and let  $t$  be the time elapsed after the explosion. Using dimensional arguments<sup>1-5</sup> or group-theoretical discussions<sup>6</sup> one can derive the following relation between  $x_s$  and  $t$ :

$$x_s = K(n, \gamma) \left( \frac{E_o}{\rho_o} \right)^{1/(2+n)} t^{2/(2+n)} \quad (2.1)$$

In Eq. (2.1),  $K(n, \gamma)$  is a proportionality factor that depends on the ratio of specific heats,  $\gamma$ , and on the dimension  $n$  of the event ( $n = 1, 2$ , or  $3$  for planar, cylindrical, or spherical symmetry, respectively),  $\rho_o$  is the density of the ambient gas, and  $E_o$  is the energy released, per unit area for a planar wave, per unit length for a cylindrical wave, and the total energy released for a spherical wave.

The proportionality factor  $K(n, \gamma)$  is determined from the requirement that the total energy contained in the blast wave must be equal to the total energy released, i.e.,

$$E_o = \frac{\pi^{n/2}}{\Gamma(\frac{n}{2})} \int_0^{x_s(t)} \rho(x, t) \left( u^2(x, t) + \frac{2}{\gamma-1} \frac{p(x, t)}{\rho(x, t)} \right) x^{n-1} dx. \quad (2.2)$$

In Eq. (2.2),  $\rho(x,t)$  is the density of the gas,  $u(x,t)$  is its particle velocity, and  $p(x,t)$  is pressure.  $\Gamma(\frac{n}{2})$  is the gamma function. It assumes the values  $\sqrt{\pi}$ , 1, and  $1/2 \sqrt{\pi}$  for  $n = 1, 2$ , and 3, respectively.

If one substitutes in Eq. (2.2) for  $\rho$ ,  $u$ , and  $p$  the self-similar solution formulas of Section 2.2. then the integral is found to be independent of time and proportional to  $E_0$ . It does, however, depend on  $n$ ,  $\gamma$  and  $K(n,\gamma)$ . Therefore, Eq. (2.2) provides for each set of  $n$  and  $\gamma$  a corresponding value for the proportionality factor  $K(n,\gamma)$ . The integral has to be evaluated numerically. We shall return to the numerical quadrature in Section 3.2.

The pressure and the particle velocity at the shock are expressed most conveniently in terms of the shock velocity  $U = dx_s(t)/dt$ . The latter is, according to Eq. (2.1)

$$U(t) = \frac{2}{2+n} \frac{x_s(t)}{t} . \quad (2.3)$$

The particle velocity immediately behind the shock is

$$u_s(t) = \frac{2}{\gamma+1} U(t) . \quad (2.4)$$

and the corresponding pressure is

$$p_s(t) = \frac{2}{\gamma+1} \rho_0 U^2(t) , \quad (2.5)$$

where  $\rho_0$  is the density of the ambient gas. The density immediately behind the shock is

$$\rho_s = \frac{\gamma+1}{\gamma-1} \rho_0 . \quad (2.6)$$

We reiterate that Eqs. (2.3) through (2.6) are "strong shock" relations, i.e., they are based on the assumption that the ambient pressure  $p_0$  is negligible compared to the shock pressure  $p_s$ .

## 2.2. Flow Field Formulas

The closed form expressions for the flow field behind a strong blast are given, e.g., in References 4, 5, 6, and 8. In this report, we formulate the solution in a somewhat simpler but algebraically equivalent form.

The formulas express the distance, particle velocity, pressure, and density in terms of time and a dimensionless parameter  $v$ . The expressions are valid for ratios of specific heat that satisfy the conditions

$$1 < \gamma < 7 \quad \text{and} \quad \gamma \neq 2. \quad (2.7)$$

The case  $\gamma = 2$  requires a special treatment because some of the formulas become singular for that value of  $\gamma$ . The solution in this special case is simpler than the general solution<sup>12</sup>, but we have not included the special solution in our computer programs.

The formulas for the flow variables are

$$x = x_s(t) \cdot y(v), \quad (2.8)$$

$$u = u_s(t) \cdot \frac{v}{v_2} \cdot y(v), \quad (2.9)$$

$$p = p_s(t) \cdot g(v), \quad (2.10)$$

$$\rho = \rho_s \cdot h(v). \quad (2.11)$$

The time functions and  $\rho_s$ , appearing in Eqs. (2.8) through (2.11) are defined in Section 2.1. The functions  $y(v)$ ,  $g(v)$ , and  $h(v)$  are defined as follows:

$$y(v) = \left( \frac{v}{v_2} \right)^{-C_1} \left( \frac{v-v_1}{v_2-v_1} \right)^{C_2} \left( \frac{1-av}{1-av_2} \right)^{C_3}, \quad (2.12)$$

$$g(v) = \left( \frac{v}{v_2} \right)^{nC_1} \left( \frac{v_1\gamma-v}{v_1\gamma-v_2} \right)^{C_4} \left( \frac{1-av}{1-av_2} \right)^{C_5}, \quad (2.13)$$

$$h(v) = \left( \frac{v-v_1}{v_2-v_1} \right)^{1-2C_2} \left( \frac{v_1\gamma-v}{v_1\gamma-v_2} \right)^{C_4-1} \left( \frac{1-av}{1-av_2} \right)^{C_5-2C_3}. \quad (2.14)$$

<sup>12</sup>A. Sakurai, "Blast Wave Theory", Mathematics Research Center Technical Summary Report 497 (September 1964).

The dimensionless parameter  $v$  varies between

$$\left. \begin{aligned} v_1 &= \frac{2}{(n+2)\gamma} \\ \text{and} \\ v_2 &= \frac{4}{(n+2)(\gamma+1)} \end{aligned} \right\} \quad (2.15)$$

The value  $v = v_1$  corresponds to the distance  $x = 0$  from the explosion, and the value  $v = v_2$  corresponds to a point on the shock.

The other constants in Eqs. (2.12) through (2.14) are

$$a = 1 + \frac{n}{2} (\gamma - 1) \quad , \quad (2.16)$$

$$\left. \begin{aligned} C_1 &= \frac{2}{2+n} \quad , \\ C_2 &= \frac{\gamma-1}{2\gamma+n-2} \quad , \\ C_3 &= \frac{(\gamma-2) 2n}{(2+n)(n\gamma-n+2)} - \frac{(2+n)(\gamma-1)\gamma}{(2\gamma+n-2)(n\gamma-n+2)} \quad , \\ C_4 &= \frac{\gamma}{\gamma-2} \quad , \\ C_5 &= \frac{(2\gamma+n-2)2n}{(2+n)(n\gamma-n+2)} + \frac{(2+n)(\gamma-1)\gamma}{(2-\gamma)(n\gamma-n+2)} \quad . \end{aligned} \right\} \quad (2.17)$$

The constant  $K(n, \gamma)$  enters the formulas (2.8), (2.9), and (2.10) through the time functions  $x_s$ ,  $u_s$ , and  $p_s$ , respectively. The constant is determined by substituting Eqs. (2.8) through (2.10) into the integral (2.2). After cancelling out the factor  $E_0$  one thereby obtains the relation

$$1 = K(n, \gamma)^{n+2} \cdot B \quad ,$$

i. e.

$$K(n, \gamma) = B^{-1/(n+2)}, \quad (2.18)$$

where

$$B = \frac{\pi^{n/2}}{\Gamma(\frac{n}{2})} \cdot \left(\frac{2}{2+n}\right)^2 \cdot \frac{4}{\gamma^2-1} \cdot \int_{v_1}^{v_2} \phi(v) dv, \quad (2.19)$$

$$\phi(v) = \left\{ D_1^{2+(2+n)C_1} D_2/D_4 + 1 \right\} D_2^{nC_2} D_3^{C_5+nC_3} D_4^{C_4} D_5, \quad (2.20)$$

$$\left. \begin{aligned} D_1 &= \frac{v}{v_2}, \\ D_2 &= \frac{v-v_1}{v_2-v_1}, \\ D_3 &= \frac{1-av}{1-av_2}, \\ D_4 &= \frac{v_1\gamma-v}{v_1\gamma-v_2}, \\ D_5 &= -\frac{C_1}{v} + \frac{C_2}{v-v_1} - \frac{C_3a}{1-av}. \end{aligned} \right\} \quad (2.21)$$

Eqs. (2.8) through (2.21) provide, together with the shock formulas of Section 2.1, explicit expressions for the flow profiles, i.e., formulas for the computation of the flow field for any fixed time  $t$ .

### 2.3. Computation of Flow Histories

Flow history, i.e., the description of the flow at a fixed position, can be obtained from the formulas of Sections 2.1 and 2.2 after a simple manipulation. Let  $X$  be the distance from the explosion at which the flow history is to be computed. The shock arrival time at  $x = X$  is, according to Eq. (2.1)

$$t_s(X) = K(n, \gamma)^{-(n+2)/2} \left( \frac{\rho_o}{E_o} \right)^{1/2} X^{(n+2)/2} \quad (2.22)$$

On the other hand, one obtains from Eqs. (2.1) and (2.8) the relation

$$t = K(n, \gamma)^{-(n+2)/2} \left( \frac{\rho_o}{E_o} \right)^{1/2} \left( \frac{X}{y(v)} \right)^{(n+2)/2} \quad (2.23)$$

Therefore,

$$t = t_s(X) y(v)^{-(n+2)/2} \quad (2.24)$$

If one varies in Eq. (2.24) the parameter  $v$  between  $v=v_2$  and  $v=v_1$ , then one obtains time values between  $t=t_s(X)$  and infinity, i.e., all times after shock arrival at  $X$ . The corresponding values of the other flow variables are obtained by substituting  $v$  and the corresponding value of  $t$  (computed using Eq. (2.24)) into the Eqs. (2.9), (2.10) and (2.11).

We notice in passing that Eq. (2.8) can be simplified for these calculations by substituting in it the definition of  $u_s(t)$  and  $v_2$ . The result is

$$u = x(t, v) \cdot \frac{v}{t}, \quad (2.25)$$

where  $x(t, v)$  is given by Eq. (2.8). For flow history computations  $x(t, v) \equiv X$ , of course.

#### 2.4. Particle Path Formulas

Closed formulas for the particle trajectories (or paths) were derived by Laporte and Chang<sup>6,7</sup>. We shall quote their formulas in a somewhat simpler but algebraically equivalent form.

The trajectory of a particle which leaves the shock at the time  $t = T$  is given by

$$\left. \begin{aligned} t &= T \cdot w(v) \\ x &= x_s(t) \cdot y(v) \end{aligned} \right\} \quad (2.26)$$

where

$$w(v) = \frac{v}{v_2} \cdot \left( \frac{\gamma v_1 - v}{\gamma v_1 - v_2} \right)^{P_1} \left( \frac{v - v_1}{v_2 - v_1} \right)^{P_2} \left( \frac{1 - av}{1 - av_2} \right)^{P_3} \quad (2.27)$$

$$\left. \begin{aligned} P_1 &= (2\gamma + n - 2)(2 + n)(1 - \gamma)\gamma \cdot P_0^{-1} \\ P_2 &= (2 + n)n(\gamma - 2)(1 - \gamma)\gamma \cdot P_0^{-1} \\ P_3 &= (2 + n)^2(\gamma - 1)^2\gamma \cdot P_0^{-1} - 1 \end{aligned} \right\} \quad (2.28)$$

$$P_0 = 2(2n\gamma + 2\gamma - n + 2)(2 - n - 2\gamma) + 2(n\gamma + 4)n(2 - \gamma)\gamma + 2(n\gamma - n + 2)(2 + n)(\gamma^2 - 1) \quad (2.29)$$

The constants  $v_1$ ,  $v_2$ , and  $a$  are defined by Eqs. (2.15) and (2.16), respectively. The function  $y(v)$  is defined by Eq. (2.12), and  $x_s(t)$  is defined by Eq. (2.1).

In Eq. (2.26), one obtains with  $v=v_2$  the initial point of the particle path, i.e., a point on the shock and with the coordinates  $t=T$  and  $x=x_s(T)$ . By decreasing the parameter  $v$  to  $v=v_1$  one obtains  $t$  and  $x$  values that approach infinity, thus covering the complete particle trajectory.

The flow variables  $u$ ,  $p$  and  $\rho$  are obtained at any point of the path line by substituting the corresponding values of  $v$  and  $t$  into the Eqs. (2.9), (2.10) and (2.11), respectively.

## 2.5. Mach-Line Formulas

Closed form expressions for the Mach-lines in a strong blast field were derived by Laporte and Chang<sup>6,7</sup>. We shall present their formulas in a somewhat simpler but algebraically equivalent form.

The Mach-lines are given in terms of the dimensionless parameter  $v$  by

$$\left. \begin{aligned} t &= M_{\pm} \cdot t_1(v) \cdot e^{\pm t_2(v)} , \\ x &= x_s(t) \cdot y(v) , \end{aligned} \right\} \quad (2.30)$$

where

$$\left. \begin{aligned} t_1(v) &= \frac{av}{1-av} , \\ t_2(v) &= \frac{1}{\sqrt{b_1+b_2}} \arcsin \frac{b_1+b_2 av_2}{av_1(1-av)} , \end{aligned} \right\} \quad (2.31)$$

$$\left. \begin{aligned} b_1 &= \left[ 2\gamma av_1 - (\gamma+1) \right] \frac{av_1}{\gamma-1} , \\ b_2 &= \left[ 2 - (\gamma+1)av_1 \right] \frac{1}{\gamma-1} . \end{aligned} \right\} \quad (2.32)$$

The function  $y(v)$  is defined by Eq. (2.12),  $x_s(t)$  is defined by Eq. (2.1), and the constants  $v_1$ ,  $v_2$  and  $a$  are defined by Eqs. (2.15) and (2.16), respectively.

The constants  $M_+$  and  $M_-$  in Eq. (2.30) correspond to the positive and negative sign of the exponent function  $t_2(v)$ , respectively, and determine the two Mach-lines passing through each point in the  $x, t$ -plane.

In order to calculate a specific characteristic one might specify as initial point either a point on the shock or a point with  $x=0$  (i.e., at the location of the explosion).

Let the initial point be on the shock at a distance  $X$  from the explosion. Then the corresponding shock arrival time  $T=t_s(X)$  can be computed with Eq. (2.22). The two constants  $M_+$  and  $M_-$  which define the two characteristics crossing at  $(X, T)$  are obtained by substituting  $v=v_2$  and  $t=T$  into Eq. (2.30). The result is



$$\left. \begin{aligned} M_+ &= \frac{T}{t_1(v_2)} e^{-t_2(v_2)} \\ M_- &= \frac{T}{t_1(v_2)} e^{+t_2(v_2)} \end{aligned} \right\} \quad (2.33)$$

If the starting point for the calculation of the characteristics is chosen at the explosion ( $x=0$ ) and time  $t=T$ , then the constants  $M_+$  and  $M_-$  are obtained from Eq. (2.30) for  $v=v_1$ . Substituting  $v=v_1$  into Eq. (2.31) one obtains

$$t_2(v_1) = - \frac{\pi}{2} / \sqrt{b_1+b_2} \quad (2.34)$$

Therefore, in this case the formulas for the constants simplify to

$$\begin{aligned} M_+ &= \frac{T}{t_1(v_1)} \exp\left(\frac{\pi}{2} / \sqrt{b_1+b_2}\right) \\ M_- &= \frac{T}{t_1(v_1)} \exp\left(-\frac{\pi}{2} / \sqrt{b_1+b_2}\right) \end{aligned} \quad (2.35)$$

Once the constants  $M_+$  and  $M_-$  are found, then all other points of each characteristic are obtained by letting the parameter  $v$  vary between  $v_1$  and  $v_2$ . The former value corresponds to a point with  $x=0$  and the latter value corresponds to a point on the shock.

The flow variables  $u$ ,  $p$ , and  $\rho$  along the Mach-lines are obtained by substituting the corresponding values of  $v$  and  $t$  into the Eqs. (2.9), (2.10), and (2.11), respectively.

### 3. COMPUTER PROGRAMS

#### 3.1. General Comments

The computer program package consists of six subroutines for flow calculations and two auxiliary routines. One of the subroutines, SBLPREP, is a preparation routine. It computes a set of constants from input data (ambient conditions and energy released) and makes the constants available to the other programs via a labeled COMMON. The preparation

routine must be called first, before any of the other flow calculation routines can be called.

The other five flow calculation routines compute, respectively, the shock, the flow profile at a given time, the flow history at a given distance, a particle path with corresponding flow variables, and a pair of Mach-lines with corresponding flow variables.

The two auxiliary routines are used by the preparation routine for its calculation of the constants.

All arguments are assumed to be expressed in SI base units.

### 3.2. Preparation Routine SBLPREP

The preparation routine SBLPREP computes the values of all constants that are needed for flow calculations and appear in the formulas of Section 2. The subroutine must be called (once) before any of the actual flow calculation routines can be used. The preparation routine is called by the statement

Call SBLPREP (N,P,TEM,GAM,AMOL,ENCHRG,NBAD)

The arguments are

N = dimension of the space (N=1,2, or 3);

P = ambient pressure (in pascals);

TEM = ambient temperature (in kelvins);

GAM = ratio of specific heats; GAM is restricted by Eq. (2.7);

AMOL = molar mass of the ambient gas (in kg/mole);

ENCHRG = energy released by the charge (in  $\text{J}\cdot\text{m}^{N-3}$ ; one kiloton TNT equivalent is  $4.184\cdot 10^{12}\text{J}$ , \*).

NBAD = an error indicator; it is set equal to zero by SBLPREP after computation of all constants; a non-zero NBAD on return indicates that the constants cannot be computed with the values given by the arguments N through ENCHRG.

\*The equivalence  $1 \text{ kton TNT} = 4.184\cdot 10^{12}\text{J}$  is given in reference 13. In reference 14, p. 13, footnote, a kton TNT is defined as  $10^{12}$  calories without specifying the calorie type, and in the same reference, pp. 13 and 647 one finds the conversion factors 4.18, 4.187, and  $4.2\cdot 10^{12}$ . According to reference 13, a calorie has a value between 4.1819J and 4.19002 J, depending on its type. In the present context, the appropriate type seems to be a thermochemical calorie" which is defined as 4.184J.

<sup>13</sup>American National Standards Institute, American Society for Testing Materials, "Standard for Metric Practice", ASTM E380-76, ASTM, Philadelphia, PA, 1976.

<sup>14</sup>Samuel Glasstone and Philip F. Dolan, eds, "The Effects of Nuclear Weapons", 3rd edition, U.S. Department of Defense and U.S. Department of Energy, 1977.

If any of the quantities P through AMOL are missing (i.e., are not positive), then the following default values are used instead of the missing quantities:

$$P = 101.325 \text{ kPa}$$

$$TEM = 293.0 \text{ K}$$

$$GAM = 1.4$$

$$AMOL = 28.96 \text{ g/mole}$$

These values correspond to a "standard" air. The use of default values is indicated by a printed message.

If the charge energy ENCHRG is not positive, or if N or GAM are not within the indicated ranges, then the calculation of the constants is not carried out. Instead, a return with NBAD  $\neq$  0 is executed, after printing of a corresponding error message. If GAM = 2.0 is specified by the calling routine, then the computation will be carried out with the default value GAM = 1.99.

The calculated constants are stored in a labeled COMMON/SBLCOM/ and made thus available to the flow calculation routines. The contents of SBLCOM is printed by the subroutine before executing a regular return.

The calculation of all but one of the constants is trivial and involves merely the evaluation of the lengthy formulas of Section 2. The computation of the constant  $K(n, \gamma)$ , defined by Eqs. (2.18) through (2.21), requires a numerical quadrature, which is carried out using a Romberg algorithm. The algorithm is implemented by the auxiliary routines SBLROMB and SBLINTE. Because the integrand  $\phi(v)$  is singular at the lower limit  $v_1$ , the integration is done in terms of the transformed variable

$$X(v) = \left( \frac{v - v_1}{v_2 - v_1} \right)^{1/(nC_2)},$$

between the limits zero and one.

The input arguments do not contain the density of the ambient air, which is needed in the formulas. The density is calculated in SBLPREP from the given arguments by the formula.

$$\rho_0 = (P \cdot AMOL) / (TEM \cdot 8.3143),$$

where 8.3143 is the universal gas constant in J/(K·mole).

The pressure and temperature instead of density were chosen as input parameters because the former quantities are more readily available when dealing with simulations of a real explosion.

### 3.3. Shock Computation Using SBLSHCK

The subroutine SBLSHCK computes the trajectory of a shock generated by a strong blast, and corresponding shock and particle velocities, shock pressure, shock density, and dynamic pressure. The computations are done for a prescribed number of distances from the explosion, equidistantly spaced between  $X_{\min}$  and  $X_{\max}$ .

The subroutine is called by the statement

```
CALL SBLSHCK (XMIN, XMAX, NRX, X, T, USHCK, P, UPART, RHO, DYNPR, NBAD)
```

The first three arguments are to be specified by the calling program. They are

XMIN, XMAX = limits of distances (in metres) from the explosion  
between which calculations should be done.

NRX = number of nodes to be calculated between XMIN and XMAX.

The next eight arguments (X through DYNPR) are one-dimensional arrays in which the calculated results will be stored:

X = array of distances of the nodes in metres,

T = array of shock arrival times at the nodes in seconds,

USHCK = array of shock velocities ( $=dx/dt$ ) in m/s,

P = array of shock pressures in pascals (P may be also interpreted as overpressure, because the ambient pressure is assumed to be negligible);

UPART = array of particle velocities at the shock in m/s,

RHO = array of densities at the shock in  $\text{kg/m}^3$ ,

DYNPR =  $0.5 \cdot \text{RHO} \cdot \text{UPART}^2$  = array of dynamic pressures in pascals.

The last argument, NBAD, is an error indicator. It is set equal to zero if all computations have been properly carried out. If NBAD is not zero at return, then the shock has not been calculated. A printed error message explains the reason for such an error return. It can be caused, e.g., by an  $\text{NRX} < 0$ , by  $\text{XMIN} > \text{XMAX}$ , etc.

The formulas for the calculation of the shock are given in Section 2.1.

### 3.4. Flow Profile Computation Using SBLPROF

The subroutine SBLPROF computes flow profiles, i.e., the flow field for fixed times. The calculations are based on formulas for the flow field of a strong blast. The formulas are given in Section 2.2.

The subroutine is called by the statement

```
CALL SBLPROF(T,XMIN,XMAX,NRX,X,P,UPART,RHO,DYNPR,NBAD)
```

The subroutine computes for the time T a total of NRX nodes located at distances between XMIN and XMAX from the explosion. The first four arguments, T through NRX, are set by the calling program. They are:

T = time (in seconds) after the explosion, for which the profile should be computed;

XMIN, XMAX = range of distances (in metres) from the explosion, between which the profile should be computed (the actual computation will be done only for positive distances and only within the blast region);

NRX = number of nodes to be computed between XMIN and XMAX.

The remaining arguments are set by the subroutine SBLPROF. The arguments X through DYNPR are one-dimensional arrays in which the subroutine will store the coordinates of the calculated nodes and the corresponding values of flow variables. The arguments are:

X = x-coordinates of the nodes (distances from the explosion in metres);

P = pressures in pascals;

UPART = particle velocities in m/s;

RHO = densities in  $\text{kg/m}^3$ ;

DYNPR = dynamic pressures in pascals

=  $0.5 * \text{RHO} * \text{UPART} ** 2$

The last argument, NBAD, is an error indicator. It is set equal to zero by the subroutine if the profile has been calculated. If the profile cannot be calculated for the specified arguments, then NBAD

assumes a non-zero value. In such a case also an error message is printed explaining the reasons for the error return.

The computed nodes are not equidistant in terms of  $X$ . Instead, the distances between the nodes are smaller in the vicinity of the shock. Other arrangements of nodes can be obtained by calling the subroutine repeatedly, e.g., once for every two nodes, specifying  $NRX=2$  and assigning corresponding values for  $XMIN$  and  $XMAX$ .

### 3.5. Flow History Computation Using SBLHIST

The subroutine SBLHIST computes flow histories at prescribed distances from the explosion. The computations are based on formulas given in Section 2.3. The subroutine is called by the statement

```
CALL SBLHIST(X,TMIN,TMAX,NRT,T,P,UPART,RHO,DYNPR,NBAD)
```

The subroutine computes NRT nodes located at the distance  $X$  from the explosion and between the times TMIN and TMAX after the explosion. The first four arguments,  $X$  through NRT, are set by the calling program. They are:

$X$  = distance (in metres) from the explosion at which the history should be computed;

TMIN,TMAX = minimum and maximum times (in seconds) after the explosion for which the computations should be done; the first node will be calculated for  $t = \max(TMIN, t_s)$ , where  $t_s$  is the shock arrival time at the distance  $X$  (Eq. (2.22));

NRT = number of nodes to be computed.

The remaining arguments are set by the subroutine. The arguments T through DYNPR are one-dimensional arrays in which the subroutine stores the computed nodal values. The arguments are:

T = array of times (in seconds) after the explosion;

P = array of pressures (in pascals);

UPART = array of particle velocities (in m/s);

RHO = array of densities (in  $\text{kg/m}^3$ );

DYNPR = array of dynamic pressures (in pascals) =  $0.5 * RHO * UPART^2$



The last argument, NBAD, is an error indicator. It is set equal to zero by the subroutine after the computations have been completed. In case the flow history cannot be computed for the given arguments, NBAD is assigned a non-zero value and a message is printed explaining the reason for the error return.

The computed nodes are approximately equidistant in time, but the distances between the nodes are smaller in the vicinity of the shock. Other node arrangements can be obtained by calling the subroutine repeatedly with proper values of TMIN and TMAX and with specifying NRT=1 or NRT=2.

### 3.6. Particle Path Computation Using SBLPATH

The subroutine SBLPATH computes particle trajectories in a strong blast flow field. The initial point of the trajectory is a point on the shock. The end point is specified by a maximum time or a maximum distance. The computations are based on formulas given in Section 2.4.

The subroutine is called by the statement

```
CALL SBLPATH(XSHCK,TMAX,XMAX,NRNOD,X,T,P,UPART,RHO,DYNPR,NBAD)
```

The first four arguments are set by the calling routine. They are:

XSHCK = distance (in metres) of the initial point of the trajectory from the explosion, i.e., the x-coordinate of a shock point;

TMAX, XMAX = coordinates of the end point of the trajectory; the path will end either at time = TMAX, or at distance = XMAX, whichever is reached first;

NRNOD = number of nodes to be computed.

The remaining arguments are set by the subroutine. The arguments X through DYNPR are one-dimensional arrays in which the subroutine stores the coordinates and flow variable values of the NRNOD computed nodes. The arguments are

X = distances (in metres) of nodes from the explosion;

T = times (in seconds) after explosion;

P = pressures (in pascals);

UPART = particle velocities =  $dx/dt$ , (in m/s);

RHO = densities (in  $kg/m^3$ );

DYNPR = dynamic pressures (in pascals) =  $0.5 \cdot \text{RHO} \cdot \text{UPART}^2$ .

The last argument, NBAD, is an error indicator. It is set equal to zero by the subroutine if the particle path has been computed. If the path cannot be computed with the data specified by the calling program then NBAD is assigned a non-zero value and a message is printed by SBLPATH explaining the reasons for the error return.

The nodes are approximately equidistant in time.

### 3.7. Mach-Line Computation Using SBLMACH

The subroutine SBLMACH computes Mach-lines in a flow field of a strong blast. The user has to specify an initial point of the Mach-lines either on the shock, or at the site of the explosion. The subroutine then establishes both Mach-lines that pass through the specified point. The actually computed nodes are located on segments of these characteristics, defined by minimum and maximum values of distance and time. The computation is based on the formulas of Section 2.5.

The subroutine is called by the statement

```
CALL SBLMACH(XZ,TZ,XMIN,XMAX,TMIN,TMAX,NRNODE,X,T,P,UPART,RHO,DYNPR,NBAD)
```

The first seven arguments, XZ through NRNODE, are set by the calling program. They are:

XZ = distance (in metres) of the initial point from the explosion  
(XZ must be zero if TZ is positive and vice versa);

TZ = time (in seconds) after explosion for the initial point; if  
TX > 0 and XZ = 0 then the initial point is assumed to be at  
the site of the explosion; if TX = 0 and XZ > 0 then the  
initial point is assumed to be on the shock at the distance  
XZ from the explosion;

XMIN,XMAX = limits between which the Mach-lines should be calculated  
TMIN,TMAX = i.e., limits of a computational "window".

NRNODE = array containing the numbers of nodes to be computed; the array  
contains two values, NRNODE(1) is the number of nodes for the  
characteristic with  $dx/dt > 0$ , and NRNODE(2) is the number of  
nodes for the other characteristic.

The remaining seven arguments are set by the subroutine. The arguments X through DYNPR are two-dimensional arrays in which the subroutine stores the nodal values of the characteristics. The dimensions of the arrays are (2,NMAX), where NMAX is the maximum number of nodes



to be computed i.e., NMAX must be larger than or equal to  $\max(\text{NRNODE}(1), \text{NRNODE}(2))$ . The nodes with the index (1, K) belong to the characteristic with  $dx/dt > 0$ , and the nodes with the index (2, K) belong to the other characteristic.

For clarity, we show the arguments with their dimensions. The arguments are:

$X(2, \text{NMAX})$  = distances (in metres) of the nodes from the explosion;

$T(2, \text{NMAX})$  = times (in seconds) after the explosion;

$P(2, \text{NMAX})$  = pressures (in pascals);

$\text{UPART}(2, \text{NMAX})$  = particle velocities (in m/s);

$\text{RHO}(2, \text{NMAX})$  = densities (in  $\text{kg/m}^3$ );

$\text{DYNPAR}(2, \text{NMAX})$  = dynamic pressures (in pascals) =  $0.5 * \text{RHO} * \text{UPART}^2$

The argument NBAD is an error indicator. It is set equal to zero by the subroutine after the Mach-lines have been computed. If they cannot be computed for the specified input arguments, then NBAD is assigned a non-zero value and an error message is printed by SBLMACH explaining the reason for the error return.

If only one Mach-line should be computed, then the user may specify  $\text{NRNODE} = 0$  for the other Mach-line.

#### 4. EXAMPLE

As an example we present calculations of a one megaton TNT explosion in air assuming spherical symmetry. The calculations were started by calling the preparation subroutine with the statement

```
CALL SBLPREP (3,0.,0.,0.,0.,4.184 E15,NBAD).
```

The first argument advises the program that the explosion is spherically symmetric ("three-dimensional"). The next four arguments describe the ambient conditions. Because they are not positive the subroutine SBLPREP assigns standard air default values for the ambient pressure, temperature, ratio of specific heats and molar mass (see Section 3.2). The charge energy is specified as  $4.184 \cdot 10^{15} \text{J}$ , which equals one megaton TNT equivalent, expressed in joules. (The conversion factor is  $4.184 \cdot 10^9 \text{J}$  for one ton TNT equivalent<sup>13</sup>.)

The calculation results are shown in Figures 1 through 6. Figures 1 and 2 show flow profiles at 0.4 seconds after the explosion. The profiles were obtained by calling the subroutine SBLPROF. Figures 3 and 4 present flow histories at 800 m from the explosion. They were obtained by plotting the results from the subroutine SBLHIST. Figure 5

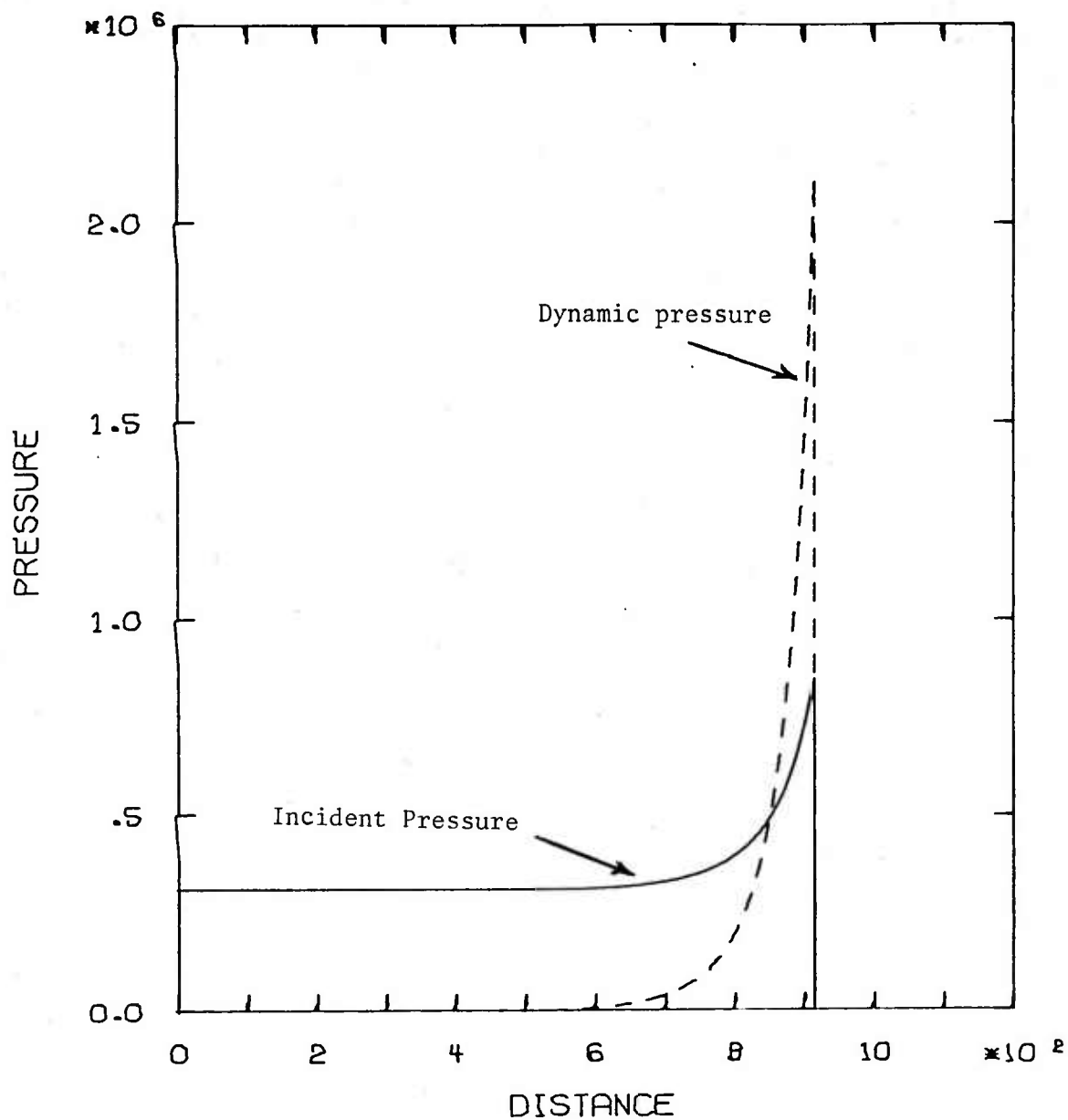


Figure 1. Pressure Profiles 0.4 Seconds After the Explosion.

Distances are expressed in metres and pressures are expressed in pascals. Dynamic pressure is defined as the product  $0.5 \cdot \rho u^2$ .

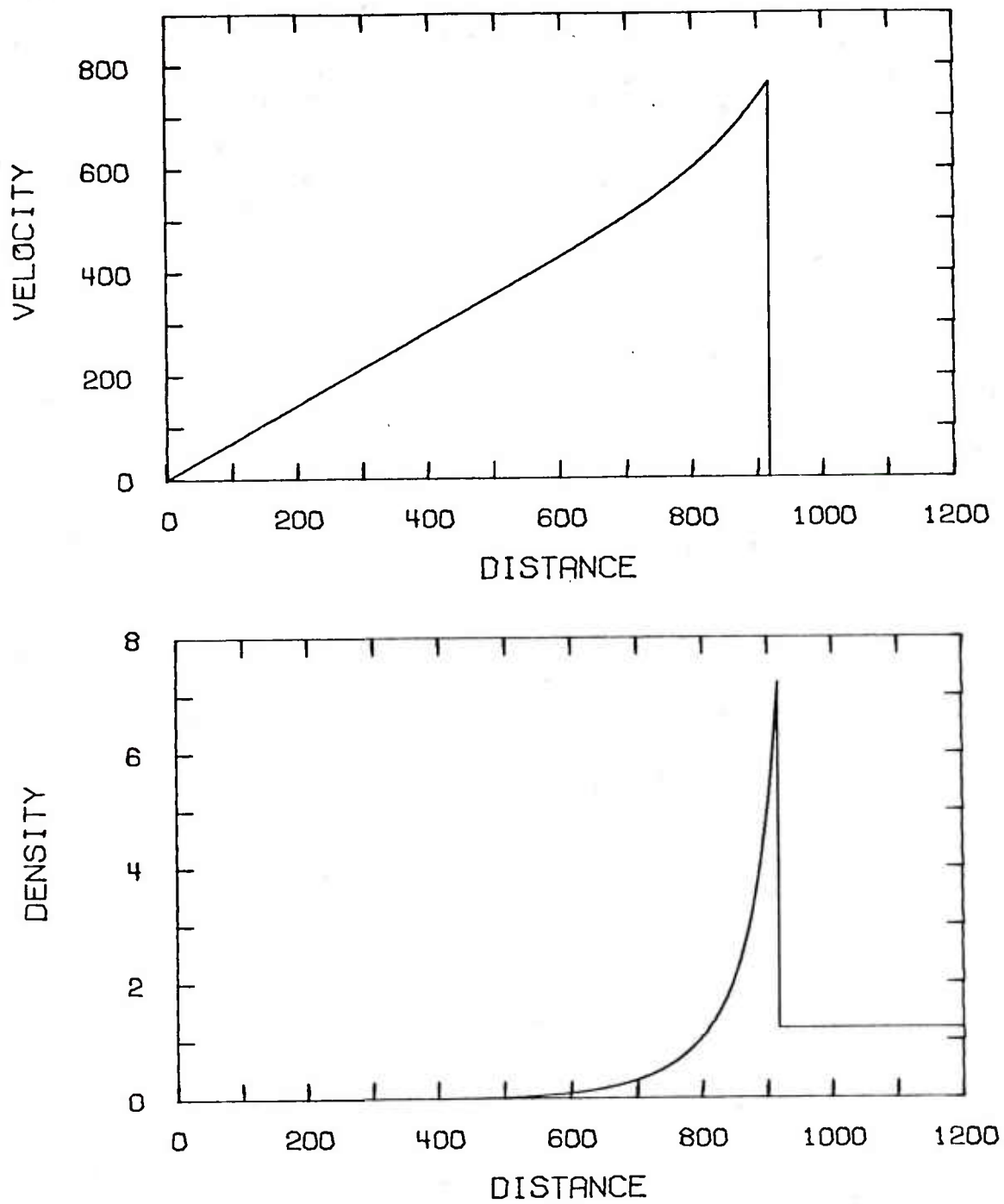


Figure 2. Velocity and Density Profiles 0.4 Seconds After the Explosion. Velocities are expressed in m/s, densities are expressed in kg/m<sup>3</sup>, and distances are expressed in metres.

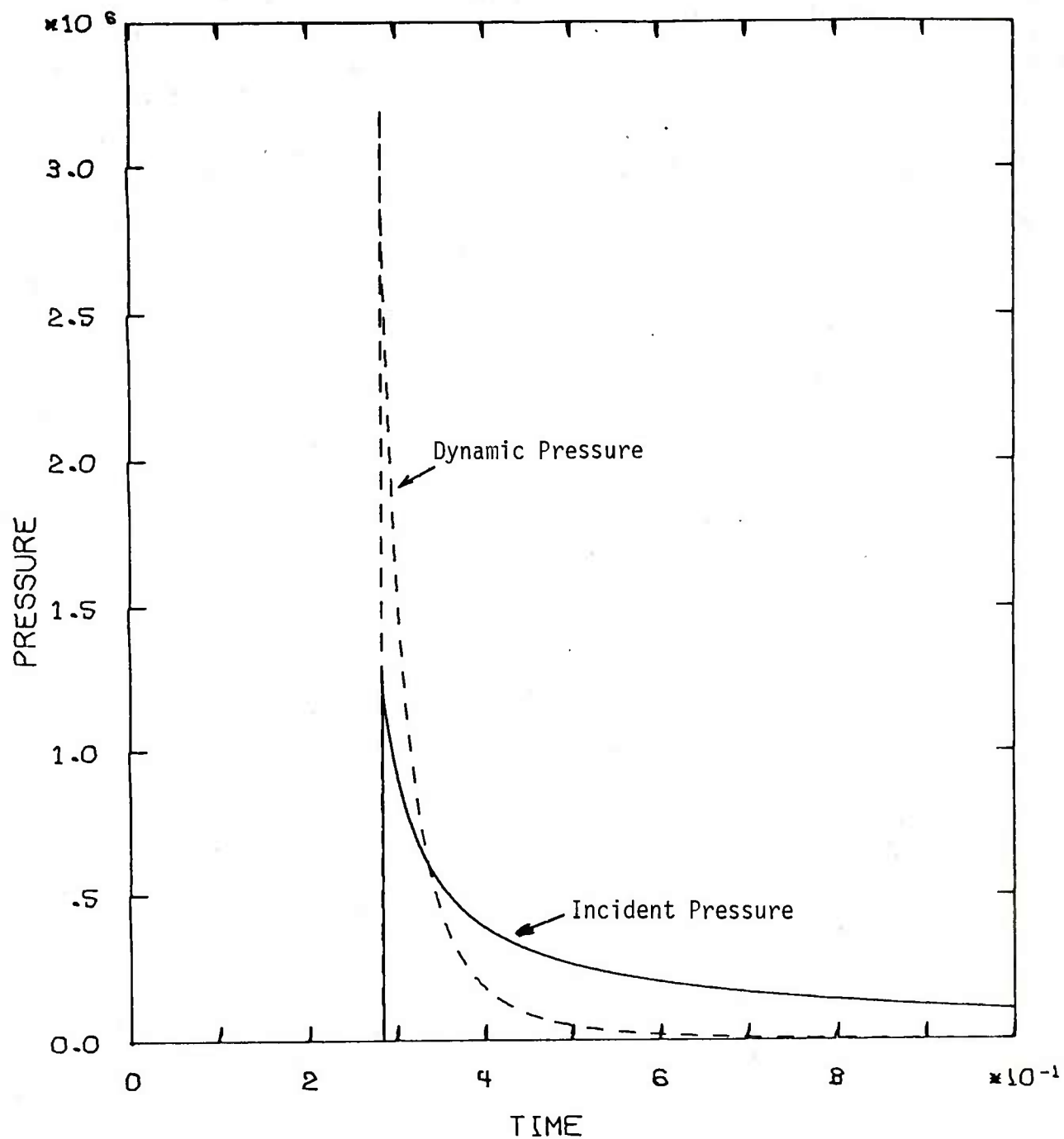


Figure 3. Pressure Histories at 800 m from the Explosion.

Times are expressed in seconds and pressures are expressed in pascals. Dynamic pressure is defined as the product  $0.5 \cdot \rho u^2$ .

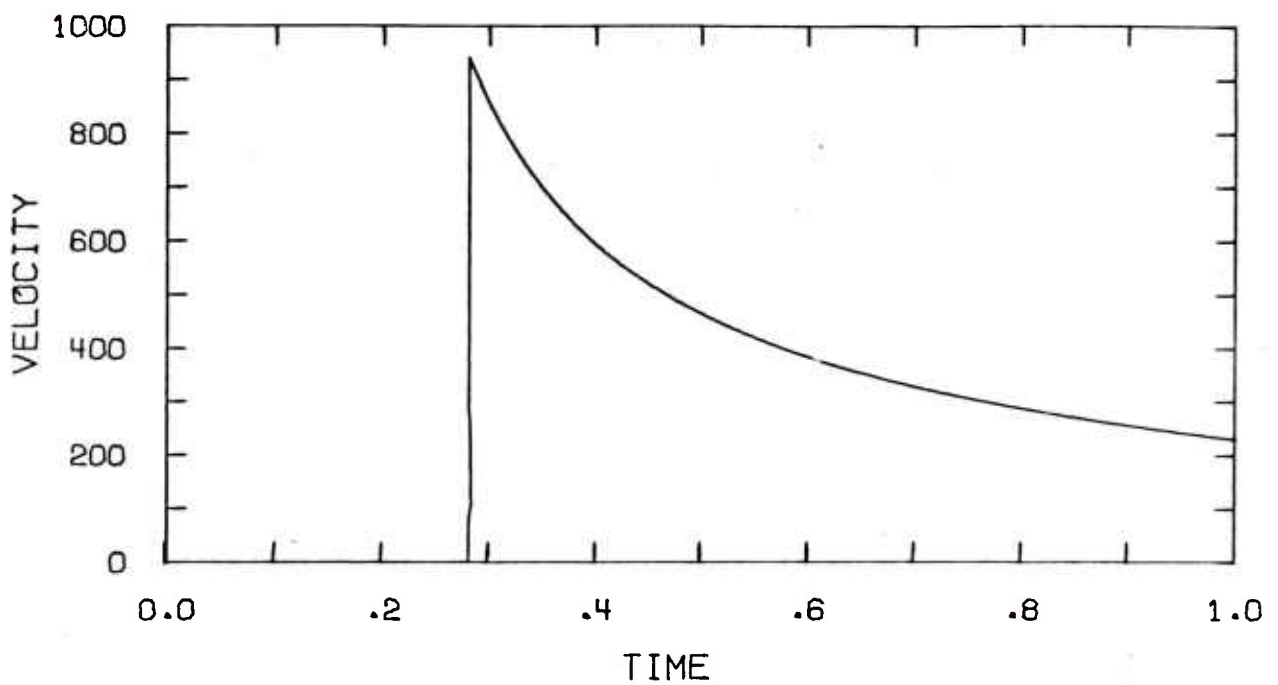
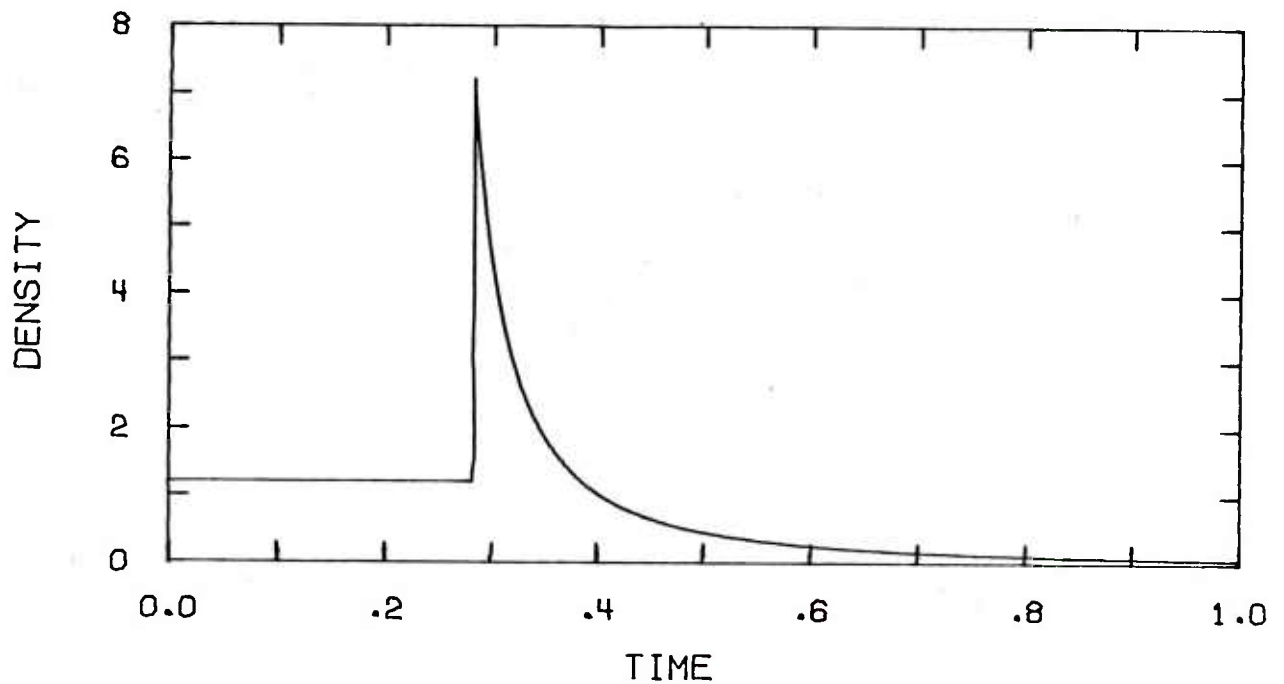


Figure 4. Velocity and Density Histories at 800 m from the Explosion.

Velocities are expressed in m/s, densities are expressed in  $\text{kg/m}^3$  and times are expressed in seconds.

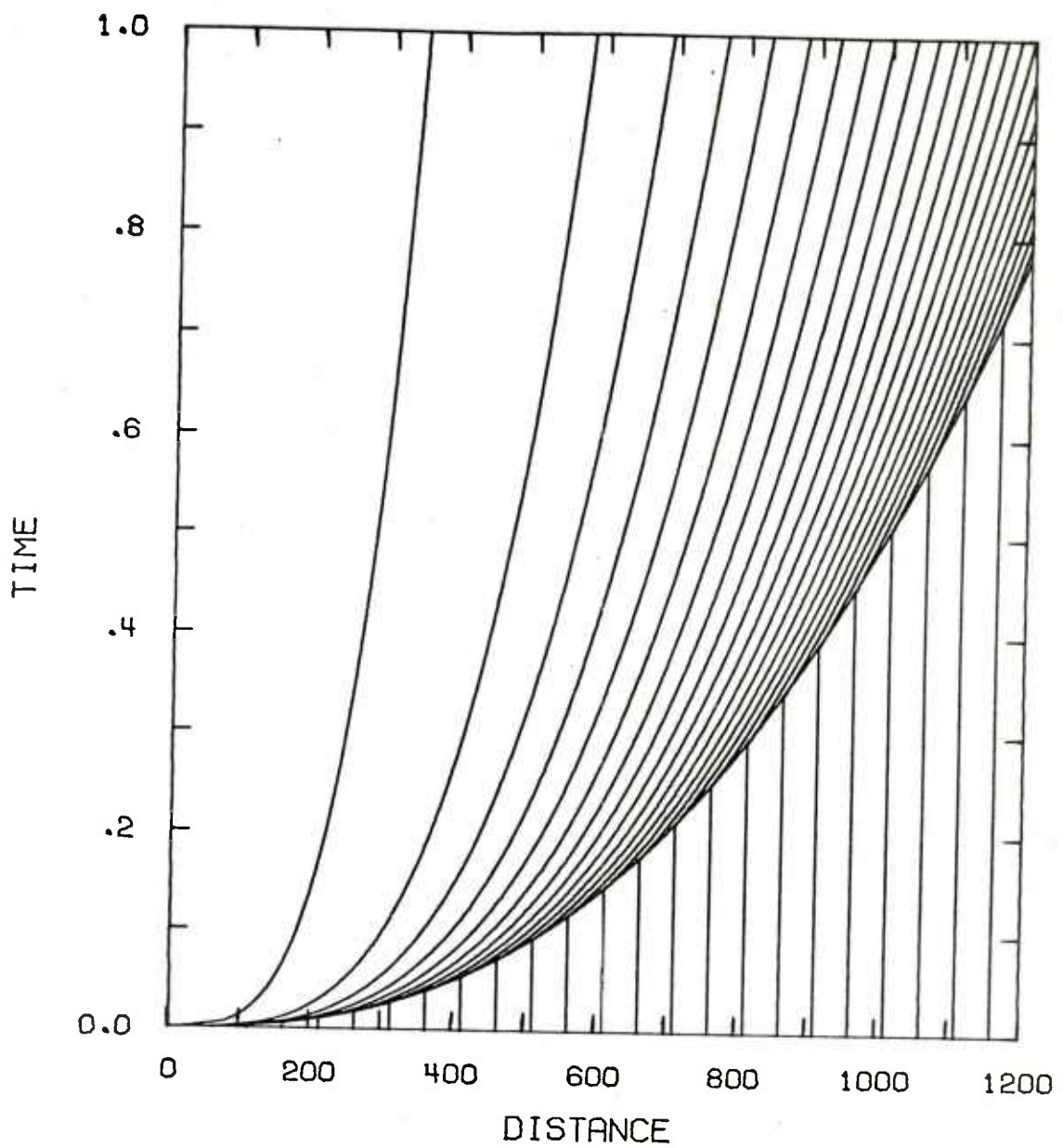


Figure 5. Particle Trajectories of a one megaton TNT Explosion, Distance is expressed in metres and time is expressed in seconds.

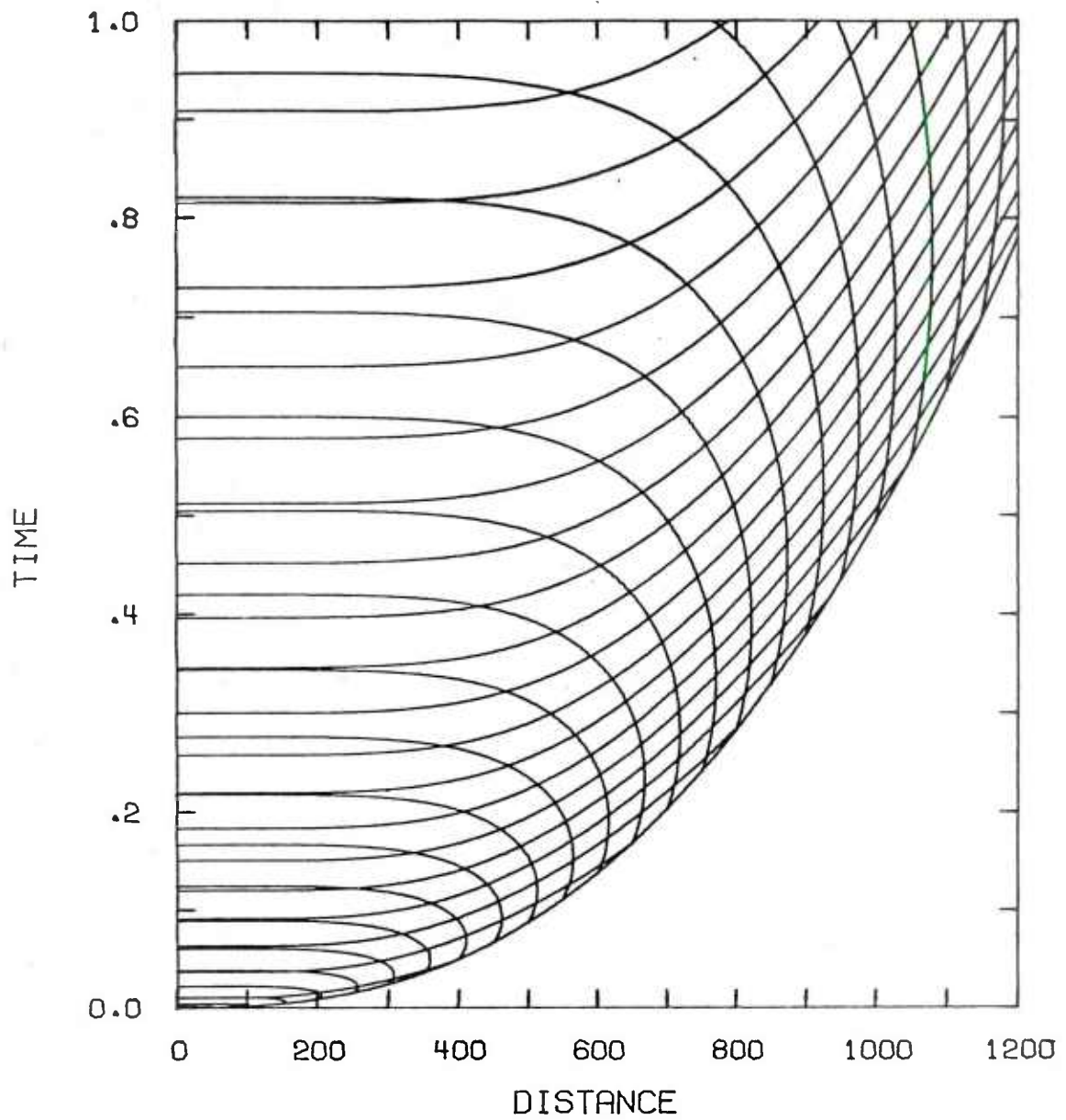


Figure 6. Mach-lines in the Blast Bubble of a one megaton TNT Explosion.

Distance is expressed in metres and time is expressed in seconds.



shows some particle paths computed using the subroutine SBLPATH, and Figure 6 shows Mach-lines computed by the subroutine SBLMACH. The shock line in Figure 5 and Figure 6 was obtained by calling the subroutine SBLSHCK.

## 5. LIMITS OF POINT SOURCE APPROXIMATIONS OF EXPLOSIONS IN AIR

The strong blast formulas of Section 2 are based on assumptions that are outlined in Section 1 and that might be satisfied at distances not too close nor too far from the explosion. In order to make the formulas useful as approximations to real explosions, one has to specify more precisely their range of applicability. In this section we shall derive quantitative conditions for which the strong blast formulas approximate spherical explosions in air. The derivations will be based on calculations in references 15 and 16. Both references report results obtained by numerical solution of so-called bursting sphere problems.

A bursting sphere problem is characterized as follows: At time  $t = 0$  one postulates an ideal gas in which an amount of internal energy  $E$  is uniformly distributed throughout a sphere of radius  $R$ . The sphere can be thought of as representing the fireball formed during the initial stages of an explosion when thermal and chemical processes dominate. The time zero is the time at which the fireball is completely formed and the hydrodynamic motion starts by forming a shock, which subsequently propagates into the ambient atmosphere. The initial velocities are assumed to be zero everywhere, and the initial density is assumed to be constant and equal  $\rho_1$  within the sphere, and constant and equal  $\rho_0$  outside the sphere. (In reference 15 it is also assumed that  $\rho_1 = \rho_0$ .) Let the corresponding initial pressures be  $p_1$  and  $p_0$ , respectively.

First we discuss the dependence of the incident shock pressure on distance, as obtained from strong blast formulas and flow field calculations, respectively. The results from bursting sphere calculations in reference 15 are displayed in Figures 7 and 8. The curves in the figures represent shock pressures in an ideal gas with the ratio of specific heats  $\gamma = 1.4$ . The ordinate in Figure 7 is the ratio  $p/p_{s0}$ , i.e., the ratio of shock pressure to initial shock pressure (at time zero). The abscissa is the ratio of distance  $r$  to the initial sphere radius  $R$ . One has the following relations between the initial shock pressure  $p_{s0}$ , the initial pressure  $p_1$  in the sphere and the energy  $E$ :

<sup>15</sup>M. Lutzky and D. Lehto, "Transformations for Scaling of Close-In Pressures from Nuclear Explosions", Naval Ordnance Laboratory Report NOLTR-66-12, Silver Spring, MD, March 1966.

<sup>16</sup>Shih Lien Huang and Pei Chi Chou, "Calculations of Expanding Shock Waves and Late-Stage Equivalence", Drexel Institute of Technology Report 125-12, Philadelphia, PA, April 1968.



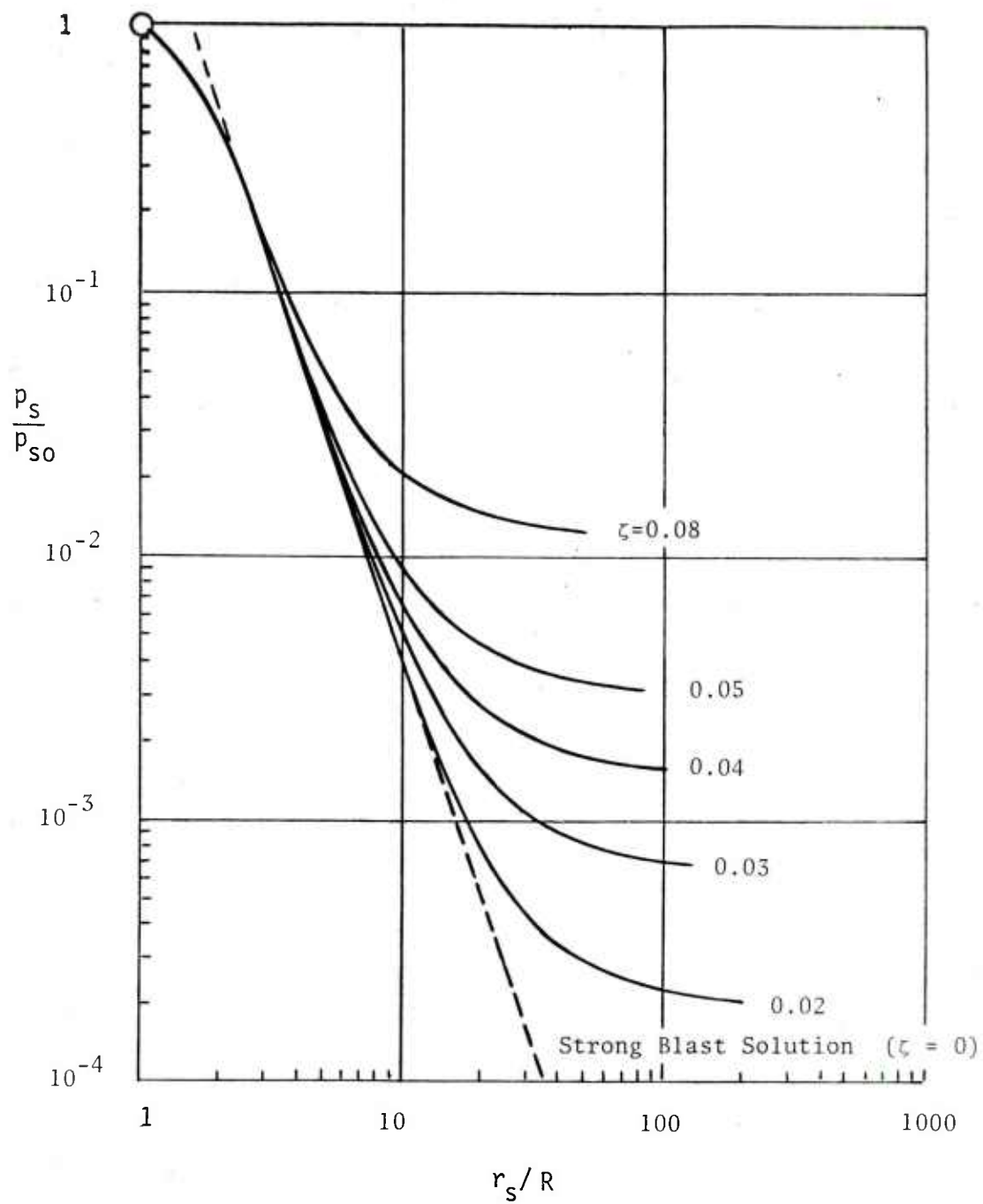


Figure 7. Shock Pressure versus Distance for Bursting Spheres.  
The curves are taken from reference 15.

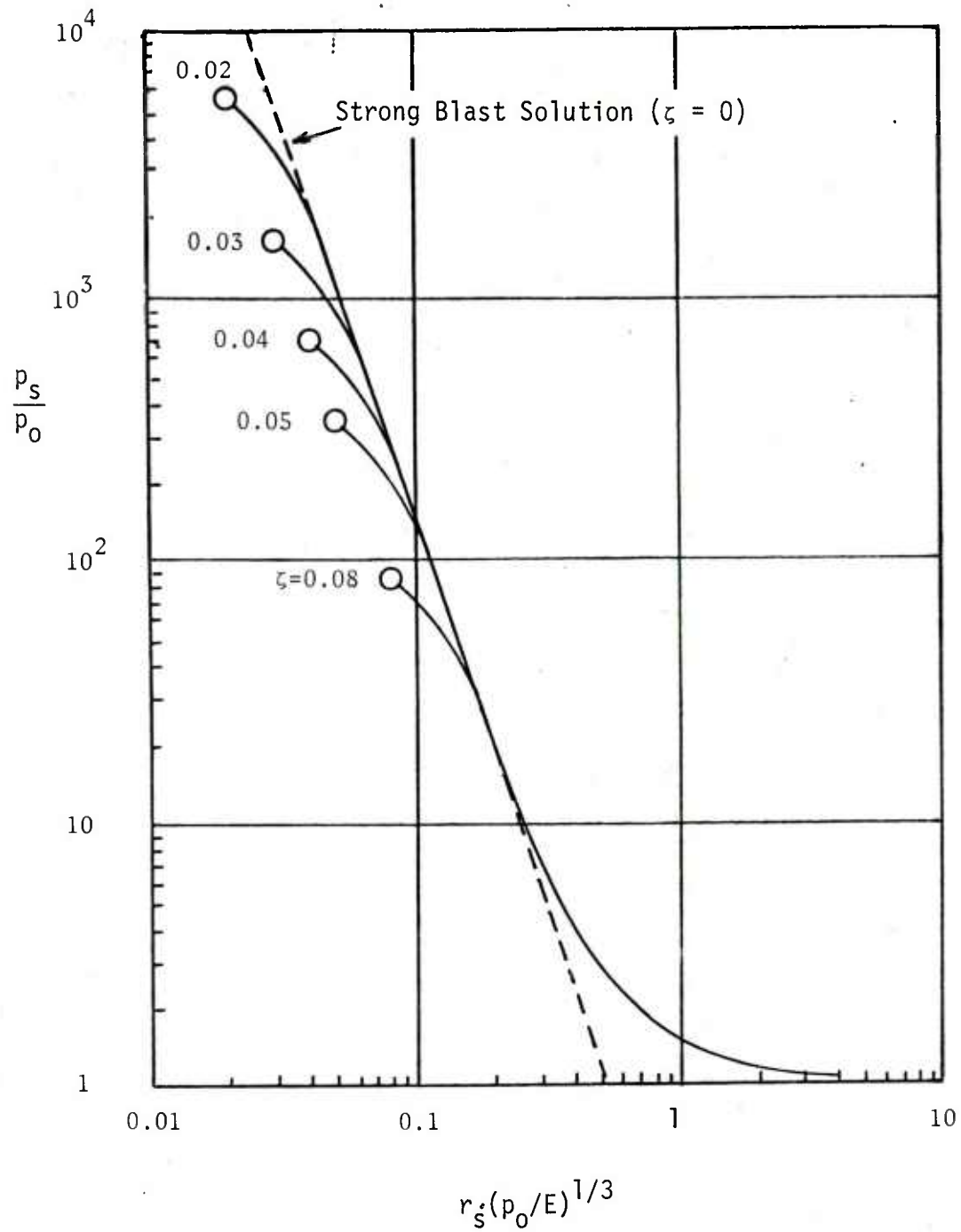


Figure 8. Shock Pressure versus Distance for Bursting Spheres (Sachs' Scales).

The figure is taken from reference 15.

$$p_{so} = 0.461 \cdot p_1 = 0.0440 \cdot E/R^3. \quad (5.1)$$

(These and other relations in this section are derived in reference 15 for  $\gamma=1.4$  and  $p_1/p_0 \gg 1$ .) The strength of the explosion is characterized by the dimensionless parameter

$$\zeta = \sqrt[3]{\frac{R^3 p_0}{E}} = 0.453 \sqrt[3]{\frac{p_0}{p_1}}. \quad (5.2)$$

The dashed line in Figure 7 represents the strong blast solution. According to Section 2.1 that solution is

$$p_s = \frac{2}{\gamma+1} \cdot \left( \frac{2}{2+n} \right)^2 \cdot K(n, \gamma)^{2+n} \cdot E / x_s^n \quad (5.3)$$

In the present case,  $n = 3$ ,  $\gamma=1.4$ ,  $K(3, 1.4) = 1.03278$  and

$$p_s = 0.1567 \cdot E/r_s^3, \quad (5.4)$$

or, by substituting Eq. (5.1) into Eq. (5.4),

$$\frac{p_s}{p_{so}} = 3.56 \cdot (R/r_s)^3 \quad (5.5)$$

It is apparent from Figure 7 that the strong blast solution approximates a finite segment of the shock curve only for events with  $\zeta < 0.08$ . In terms of initial pressures or energy, this condition can be formulated by either of the following equivalent relations

$$\left. \begin{aligned} p_{so}/p_0 &> 86, \\ p_1/p_0 &> 190, \\ E/(R^3 p_0) &> 1950. \end{aligned} \right\} \quad (5.6)$$

If the event satisfies the conditions (5.6) then the strong blast solution can be used as an approximation between  $r/R = 2.5$  and an upper limit which depends on the strength of the explosion.

The upper limit for the approximation can be obtained from Figure 8 which contains the same curves as Figure 7, but uses Sachs' scaling for its coordinates. Figure 8 shows that an upper limit is given by the condition  $p_s/p_o > 10$ .

Both conditions can be expressed either in terms of distances or in terms of pressures. In terms of distances one obtains the following interval in which the strong blast solution can be used as an approximation:

$$2.5 < \frac{r}{R} < \left\{ \begin{array}{l} 0.25/\zeta \\ 0.25 \sqrt[3]{E/(R^3 p_o)} \\ 0.547 \sqrt[3]{p_1/p_o} \\ 0.707 \sqrt[3]{p_{so}/p_o} \end{array} \right. \quad (5.7)$$

In terms of pressures one has the equivalent conditions

$$\left. \begin{array}{l} 0.0101/\zeta^3 \\ 0.0101 \cdot E/(R^3 p_o) \\ 0.106 \cdot p_1/p_o \\ 0.230 \cdot p_{so}/p_o \end{array} \right\} > \frac{p_s}{p_o} > 10. \quad (5.8)$$

In order to check either of the conditions one needs in addition to the energy  $E$  also an estimate of the fireball radius  $R$ . Such an estimate is often not available. It is, therefore, desirable to have conditions which are independent of  $a$ . If one is content with more restrictive conditions than (5.7) or (5.8), then one can replace  $p_{so}$  in the last Eq. (5.8) by an observed maximum  $p_s$  value. Let that value be  $p_{smax}$ . Then a condition for the applicability of strong blast shock formulas to explosions in air is the following set:

$$\left. \begin{aligned}
 &\frac{p_s}{p_{smax}} < 0.23, \\
 \text{and} \\
 &\frac{p_s}{p_o} > 10.
 \end{aligned} \right\} \quad (5.9)$$

An observation of  $p_{smax}$  can be also used to obtain an upper limit for the fireball radius. That limit is, according to Eq. (5.1) given by

$$R < 0.353 \sqrt[3]{E/p_{smax}} \quad (5.10)$$

The flow field generated by a bursting sphere contains in addition to the leading shock also a second shock and a contact discontinuity. Therefore, the strong blast formulas may approximate the field only within a strip in the  $r,t$ -plane behind the leading shock, where the flow is free from these singularities. In order to establish limits for approximations within the strip one needs examples of flows generated by explosions for a reasonable range of the parameter  $\zeta$ . Unfortunately such calculations have not been published. Reference 16 provides curves for one typical example only. That example has the explosion strength  $\zeta = 0.076$ , i.e., it is a limiting case where the strong blast approximation of the shock strength is not valid. (The calculation is done for the density ratio  $\rho_1/\rho_o = 1.16$ , instead of 1.0, but this is probably not significant.) One can hypothesize that the approximation by strong blast formulas will be better for more powerful explosions than in this example.

A comparison between numerical flow field calculations and strong blast formulas is done in Reference 16 in terms of relative pressure profiles  $p/p_s$ , relative particle velocity profiles  $u/u_s$  and relative sound speed profiles  $c/c_s$ . The calculations and comparisons are presented for times corresponding to strong blast shock positions between  $r/R = 3.58$  and  $r/R = 7.60$ . From Figure 7 one can see that for these positions the computed shock strength differs significantly from the strong blast shock strength. (See curve  $\zeta = 0.08$  in Figure 7). Nevertheless, the corresponding relative pressure profiles differ by no more than 20%, the strong blast formulas providing higher values. (The comparison is made within a strip behind the leading shock covering 15% of the blast bubble radius). The differences become smaller for later times. At  $r/R = 7.60$  the maximum difference is only about 1%.

The relative particle velocity profiles behave differently. Again, the strong blast formulas provide higher values. The difference is up

to 13% and the approximation does not improve as time increases.

The relative sound speed profiles are up to 10% below the strong blast solution, and the difference does not decrease as the time increases.

The gas density is proportional to  $c^{-2}$  and, therefore, the calculated density would be at large distances up to 20% higher than predicted by strong blast formulas. Consequently the calculated dynamic pressure ( $\sim u^2 c^{-2}$ ) can be expected to be about 6% lower than given by the strong blast formulas, because the deviations in sound speed and particle velocity compensate each other.

In summary, in the example, presented in Reference 16, the relative incident pressure profiles and the relative dynamic pressure profiles can be reasonably approximated by corresponding strong blast profiles, particularly at large distances. This finding is encouraging because the example represents a case in which the initial shock pressure differs significantly from the corresponding strong blast pressure. One can expect better approximations for stronger explosions (smaller parameter  $\zeta$ ), but this hypothesis should be tested by sample calculations. Such calculations can be carried out, e.g., by using the computer program described in Reference 17.

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# LIST OF SYMBOLS

$a$	constant defined by Eq. (2.16)
$c$	sound speed (m/s)
$C_1, \dots, C_5$	constants defined by Eq. (2.17)
$E$	energy contained in a bursting sphere (J)
$E_0$	energy released by the explosion ( $J \cdot m^{n-3}$ )
$g(v)$	dimensionless function defined by Eq. (2.13)
$k(v)$	dimensionless function defined by Eq. (2.14)
$K(n, \gamma)$	proportionality factor, see Eq. (2.1) and (2.18)
$M_+, M_-$	factors in Mach line formulas (2.30)
$n$	dimension of the event ( $n=1,2,3$ for planar, cylindrical, spherical events, respectively)
$p$	pressure (Pa)
$p_0$	ambient pressure (Pa)
$p_1$	initial pressure in a bursting sphere (Pa)
$p_s$	shock pressure (Pa)
$p_{s0}$	initial shock pressure in a bursting sphere event (Pa)
$p_{smax}$	maximum shock pressure observed (Pa)
$r$	distance from the center of a bursting sphere
$R$	fireball radius (m)
$t, T$	time after the explosion (s)
$t_1(v), t_2(v)$	dimensionless functions defined by Eqs. (2.13) and (2.32)
$t_s$	shock arrival time (s)
$u$	particle velocity (m/s)
$u_s$	particle velocity immediately behind the shock (m/s)
$U$	shock velocity (m/s)

# LIST OF SYMBOLS (Continued)

$v$	running parameter (see Eq. (2.15))
$v_1, v_2$	particular values of $v$ , defined by Eq. (2.15)
$w(v)$	dimensionless function defined by Eq. (2.27)
$x, X$	distance from the center of the explosion (m)
$x_s, X_s$	shock distance from the center of the explosion (m)
$y(v)$	dimensionless function defined by Eq. (2.12)
$\gamma$	ratio of specific heats
$\zeta$	explosion strength parameter defined by Eq. (5.2)
$\rho$	gas density ( $\text{kg/m}^3$ )
$\rho_0$	ambient gas density ( $\text{kg/m}^3$ )
$\rho_1$	initial gas density in a bursting sphere ( $\text{kg/m}^3$ )
$\rho_s$	gas density immediately behind the shock ( $\text{kg/m}^3$ )

APPENDIX  
LISTING OF COMPUTER PROGRAMS

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C   SUBROUTINE SBLPREP(N, AIRP, AIRT, AIRC, AIRM, CHAREN, NBAD)
    THIS IS PREPARATION ROUTINE FOR STRONG BLAST COMPUTATIONS
    COMMON/SBLCOM/NC, GAMMA, RO, CHEN, V(2), A, B(2), C(5), P(3), AK, NGOOD
    EXTERNAL SBLINTE
    DATA(NFIRST=0)
    IF(NFIRST.EQ.0) NGOOD=0
    NFIRST=1
    NBAD=0
    IDEF=0
    PRES=AIRP
    IF(PRES.GT.0.)GOTO 15
    PRES=101325.1
    IDEF=IDEF+1
15  TEMP=AIRT
    IF(TEMP.GT.0.)GOTO 25
    TEMP=293.0
    IDEF=IDEF+1
25  GAM=AIRC
    IF(GAM.GT.0.)GOTO 35
    GAM=1.4
    IDEF=IDEF+1
35  IF(GAM.NE.2.0)GOTO 38
    GAM=1.99
    IDEF=IDEF+1
38  AMOL=AIRM
    IF(AMOL.GT.0.)GOTO 65
    AMOL=0.02896
    IDEF=IDEF+1
65  IF(CHAREN.GT.0.)GOTO 85
    NBAD=1
    PRINT 75, NBAD
    RETURN
75  FORMAT(1H0,10X,30HRETURN FROM SBLPREP WITH NBAD=,I2,
    A39H BECAUSE CHARGE ENERGY IS NOT SPECIFIED)
85  IF(N.GE.1.AND.N.LE.3)GOTO 101
    NBAD=2
    PRINT 95, NBAD, N
    RETURN
95  FORMAT(1H ,10X,30HRETURN FROM SBLPREP WITH NBAD=,I2,
    A11H BECAUSE N=,I3,24H IS OUTSIDE RANGE 1 TO 3)
101 IF(GAM.GT.1.0.AND.GAM.LT.7.0)GOTO 105
    NBAD=3
    PRINT 103, NBAD, GAM
    RETURN
103 FORMAT(1H0,10X,30HRETURN FROM SBLPREP WITH NBAD=,I2,
    A 15H BECAUSE GAMMA=,1PE12.5,32H IS OUTSIDE THE RANGE 1.0 TO 7.0)
105 NC=N
    AN=N
    GAMMA=GAM
    RO=PRES*AMOL/(8.3143*TEMP)
    CHEN=CHAREN
    V(1)=2.0/((2.0+AN)*GAMMA)
    V(2)=4.0/((2.0+AN)*(1.0+GAMMA))
    A=1.0+AN*(GAMMA-1.0)*0.5
    C(1)=2.0/(2.0+AN)
    C(2)=(GAMMA-1.0)/(2.0+AN)
    C(3)=2.0*AN*(GAMMA-2.0)/((2.0+AN)*(AN+2.0))-GAMMA*

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**** 1
SBLPRE 2
SBLPRE 3
SBLPRE 4
SBLPRE 5
SBLPRE 6
SBLPRE 7
SBLPRE 8
SBLPRE 9
SBLPRE10
SBLPRE11
SBLPRE12
SBLPRE13
SBLPRE14
SBLPRE15
SBLPRE16
SBLPRE17
SBLPRE18
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SBLPRE40
SBLPRE41
SBLPRE42
SBLPRE43
SBLPRE44
SBLPRE45
SBLPRE46
SBLPRE47
SBLPRE48
SBLPRE49
SBLPRE50
SBLPRE51
SBLPRE52
SBLPRE53
SBLPRE54
SBLPRE55
SBLPRE56
SBLPRE57

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A	(2.0+AN)*(GAMMA-1.0)/((2.0*GAMMA+AN-2.0)*(AN*GAMMA-AN+2.0))	SBLPRE5
C	(4)=GAMMA/(GAMMA-2.0)	SBLPRE5
C	GAMMA=2.0 IS REPLACED BY 1.99 AT STATEMENT 35	SBLPRE6
	C(5)=AN*(2.0*GAMMA+AN-2.0)/((2.0+AN)*(AN*GAMMA-AN+2.0))	SBLPRE6
	A+(2.0+AN)*GAMMA*(GAMMA-1.0)/(2.0*(2.0-GAMMA)*(AN*GAMMA-AN+2.0))	SBLPRE62
	B(1)=(-1.0-GAMMA+2.0*GAMMA*A*V(1))*A*V(1)/(GAMMA-1.0)	SBLPRE63
	B(2)=(2.0-(GAMMA+1.0)*A*V(1))/(GAMMA-1.0)	SBLPRE64
	DA=(2.0+AN)*(AN*GAMMA**2-2.0*AN*GAMMA+AN+2.0*GAMMA-2.0)	SBLPRE65
	DB=2.0*(2.0*AN*GAMMA+2.0*GAMMA-AN+2.0)*(2.0-AN-2.0*GAMMA)	SBLPRE66
	A+2.0*AN*GAMMA*(2.0-GAMMA)*(AN*GAMMA+4.0)+2.0*(GAMMA+1.0)*DA	SBLPRE67
	P(2)=AN*GAMMA*(2.0+AN)*(2.0-GAMMA)*(GAMMA-1.0)/DB	SBLPRE68
	P(1)=P(2)*(2.0*GAMMA+AN-2.0)/(AN*(GAMMA-2.0))	SBLPRE69
	P(3)=P(2)*DA/(AN*(2.0-GAMMA)*(AN*GAMMA-AN+2.0))-1.0	SBLPRE70
C	NEXT START QUADRATURE TO COMPUTE AK.	SBLPRE71
	CALL SBLROMB(SBLINTE,0.0,1.0,SF,NBAD)	SBLPRE72
	IF(NBAD.EQ.0)GOTO 135	SBLPRE73
	PRINT 125,NBAD	SBLPRE74
	RETURN	SBLPRE75
125	FORMAT(1H0,10X,30HRETURN FROM SBLPREP WITH NBAD=,I4,	SBLPRE76
	A28H BECAUSE OF ERROR IN SBLROMB)	SBLPRE77
135	IF(N.EQ.1)FACT=1.0	SBLPRE78
	IF(N.EQ.2)FACT=3.14159265359	SBLPRE79
	IF(N.EQ.3)FACT=6.28318530718	SBLPRE80
	DAK=FACT*(2.0/(2.0+AN))**2*(4.0/(GAMMA**2-1))*SF/(AN*C(2))	SBLPRE81
	AK=1.0/DAK**((1.0/(2.0+AN)))	SBLPRE82
	NGOOD=1	SBLPRE83
	PRINT145,NC,GAM,RO,CHEN,V(1),V(2),A,B(1),B(2),(C(J),J=1,5),P(1),	SBLPRE84
	AP(2),P(3),AK,NGOOD	SBLPRE85
145	FORMAT(1H1,10X,26HCONTENTS OF COMMON/SBLCOM/,	SBLPRE86
	X30H FOR STRONG BLAST CALCULATIONS,/,1H,10X,3HNC=,	SBLPRE87
	A12,5X,4HGAM=,0PF8.5,5X,4HRHO=,1PE12.5,5X,5HCHEN=,1PE12.5,/,	SBLPRE88
	Z 1H,10X,5HV(J)=,2(2X,1PE12.5),/,	SBLPRE89
	B 1H,10X,2HA=,1PE12.5,/,1H,10X,5HB(J)=,2(2X,1PE12.5),/,	SBLPRE90
	C 1H,10X,5HC(J)=,5(2X,1PE12.5)	SBLPRE91
	D//,1H,10X,5HP(J)=,3(2X,1PE12.5),/,	SBLPRE92
	E 1H,10X,3HAK=,1PE12.5,/,1H,10X,6HNGOOD=,I2,/,	SBLPRE93
	IF(IDEF.EQ.0)GOTO 165	SBLPRE94
	PRINT155,PRES,TEMP,GAM,AMOL	SBLPRE95
155	FORMAT(1H0,10X,41HSOME OF THE FOLLOWING ARE DEFAULT VALUES,/,1H,	SBLPRE96
	A10X,44HASSIGNED BY SBLPREP AND CORRESPONDING TO AIR,/,1H,10X,	SBLPRE97
	C13HPRESSURE =,1PE12.5,3H PA,/,	SBLPRE98
	D1H,10X,13HTEMPERATURE =,1PE12.5,2H K,/,	SBLPRE99
	D1H,10X,13HGAMMA =,1PE12.5,/,	SBLPR100
	L1H,10X,13HMOLAR MASS =,1PE12.5,8H K.G/MOLE)	SBLPR101
165	RETURN	SBLPR102
	END	SBLPR103

	SUBROUTINE SBLROMB(F,A,B,FINT,NBAD)	**** 1
C	ROMBERG INTEGRATION SUBROUTINE	SBLROM 2
	DIMENSION T(15,20)	SBLROM 3
	NBAD=0	SBLROM 4
	CALL F(A,FA,NBAD)	SBLROM 5
	IF(NBAD.NE.0)RETURN	SBLROM 6
	CALL F(B,FB,NBAD)	SBLROM 7
	IF(NBAD.NE.0)RETURN	SBLROM 8
	T(1,1)=(FA+FB)*0.5	SBLROM 9
	KM=1	SBLROM10
	KMA=1	SBLROM11
15	DEN=FLOAT(KMA)*2.	SBLROM12
	FM=0	SBLROM13
	DO 25 KA=1,KMA	SBLROM14
	AC=FLOAT(1+2*(KMA-KA))/DEN	SBLROM15
	BC=FLOAT(2*KMA-1)/DEN	SBLROM16
	ARG=AC*A+BC*B	SBLROM17
	IF(B-A) 17,18,19	SBLROM18
17	ARG=AMAX1(ARG,B)	SBLROM19
	ARG=AMIN1(ARG,A)	SBLROM20
	GOTO 20	SBLROM21
18	FINT=0.	SBLROM22
	RETURN	SBLROM23
19	ARG=AMAX1(ARG,A)	SBLROM24
	ARG=AMIN1(ARG,B)	SBLROM25
20	CALL F(ARG,FH,NBAD)	SBLROM26
	IF(NBAD.NE.0)RETURN	SBLROM27
	FM=FM+FH	SBLROM28
25	CONTINUE	SBLROM29
	FM=FM/FLOAT(KMA)	SBLROM30
	T(1,KM+1)=(T(1,KM)+FM)*0.5	SBLROM31
C	THIS IS TRAPEZ. NOW COMPUTE ROMBERG	SBLROM32
	KM=KM+1	SBLROM33
	KC=1	SBLROM34
	DDEN=1.	SBLROM35
35	KC=KC+1	SBLROM36
	DDEN=DDEN*4.	SBLROM37
	T(KC,KM)=T(KC-1,KM)+(T(KC-1,KM)-T(KC-1,KM-1))/(DDEN-1.)	SBLROM38
	IF(KC.LT.KM.AND.KC.LT.15)GOTO 35	SBLROM39
	IF(KC.GE.4)GOTO 45	SBLROM40
	KMA=KMA*2	SBLROM41
	GOTO 15	SBLROM42
45	KAF=KC-3	SBLROM43
	ITEST=0	SBLROM44
	DO 55 KA=KAF,KC	SBLROM45
C	NOW TEST CONVERGENCE	SBLROM46
	TEST=ABS(T(KA,KM)-T(KA-1,KM-1))	SBLROM47
	IF(TEST.GT.ABS(T(KC,KM))*1.E-5 ) ITEST=1	SBLROM48
	IF(TEST.LE.1.E-100)GOTO 65	SBLROM49
55	CONTINUE	SBLROM50
	IF(ITEST.EQ.0) GOTO 65	SBLROM51
	IF(KM.GE.20)GOTO 65	SBLROM52
	KMA=KMA*2	SBLROM53
	GOTO 15	SBLROM54
65	FINT=T(KC,KM)*(B-A)	SBLROM55
	RETURN	SBLROM56
	END	SBLROM57

	SUBROUTINE SBLINTE(W,F,NBAD)	***** 1
C	THIS IS THE INTEGRAND ROUTINE FOR THE COMPUTATION OF AK	SBLINT 2
C	THE INTEGRAND DOES NOT CONTAIN THE PI-FACTOR ETC.	SBLINT 3
	COMMON/SBLCOM/N,GAM,RO,CHEN,V(2),A,B(2),C(5),P(3),AK,NGOOD	SBLINT 4
	NBAD=0	SBLINT 5
	EX1=2.*C(1)*FLOAT(2+N)	SBLINT 6
	EX2=C(2)*FLOAT(N)	SBLINT 7
	EX3=C(5)*2.*C(3)*FLOAT(N)	SBLINT 8
	EX4=C(4)	SBLINT 9
	IF(W-1.0) 12,27,15	SBLINT10
12	IF(W) 15,25,35	SBLINT11
15	NBAD=1	SBLINT12
	RETURN	SBLINT13
25	U=V(1)	SBLINT14
	GOTO 37	SBLINT15
27	U=V(2)	SBLINT16
	GOTO 37	SBLINT17
35	U=V(1)+(V(2)-V(1))*W**(1./EX2)	SBLINT18
37	D1=U/V(2)	SBLINT19
	D2=(U-V(1))/(V(2)-V(1))	SBLINT20
	D3=(1.0-A*U)/(1.0-A*V(2))	SBLINT21
	D4=(V(1)*GAM-U)/(V(1)*GAM-V(2))	SBLINT22
	D5=C(2)-(U-V(1))*(C(1)/U+C(3)*A/(1.0-A*U))	SBLINT23
	F=(D1**EX1*D2/D4+1.0)*D3**EX3*D4**EX4*D5	SBLINT24
	RETURN	SBLINT25
	END	SBLINT26



C	SUBROUTINE SBLSHCK(RMIN,RMAX,NR,R,T,US,P,UP,RO,DP,NBAD)	****	1
C	THIS COMPUTES STRONG SHOCK BETWEEN RMIN AND RMAX	SBLSHC	2
C	NR = NUMBER OF NODES TO BE COMPUTED	SBLSHC	3
C	R = DISTANCE	SBLSHC	4
C	T = SHOCK ARRIVAL TIME	SBLSHC	5
C	US = SHOCK VELOCITY	SBLSHC	6
C	P = INCIDENT SHOCK PRESSURE	SBLSHC	7
C	UP = PARTICLE VELOCITY	SBLSHC	8
C	RO = DENSITY	SBLSHC	9
C	DP = DYNAMIC PRESSURE = $0.5*RO*UP**2$	SBLSHC	10
C	NBAD = ERROR INDICATED BY NBAD,NE,0	SBLSHC	11
	DIMENSION R(1),T(1),US(1),P(1),UP(1),RO(1),DP(1)	SBLSHC	12
	COMMON/SBLCON/NC,GAM,ROC,CHEN,V(2),A,B(2),C(5),PC(3),AK,NGOOD	SBLSHC	13
	IF(NGOOD,NE,0) GOTO 10	SBLSHC	14
	NBAD=11	SBLSHC	15
	PRINT 11,NBAD	SBLSHC	16
	RETURN	SBLSHC	17
11	FORMAT(1H0,10X,30HRETURN FROM SBLSHCK WITH NBAD=,I3,	SBLSHC	18
	A26H AND WITHOUT COMPUTATIONS,6/,1H,10X,	SBLSHC	19
	B34HSUBROUTINE SBLPREP MUST BE CALLED ,	SBLSHC	20
	C39H BEFORE OTHER SBL-ROUTINES CAN BE USED,/,)	SBLSHC	21
10	NBAD=0	SBLSHC	22
	IF(NR,GE,1)GOTO 25	SBLSHC	23
	NBAD=1	SBLSHC	24
	PRINT 15,NBAD,NR	SBLSHC	25
	RETURN	SBLSHC	26
15	FORMAT(1H0,10X,30HRETURN FROM SBLSHCK WITH NBAD=,I2,	SBLSHC	27
	A12H BECAUSE NR=,I4)	SBLSHC	28
25	R(1)=AMAX1(RMIN,0.)	SBLSHC	29
	ANT=2+NC	SBLSHC	30
	EX=ANT/2.0	SBLSHC	31
	TFACT=SQRT(ROC/CHEN)/AK**EX	SBLSHC	32
	ROFACT=((GAM+1.0)/(GAM-1.0))*ROC	SBLSHC	33
	DO 55 KA=1,NR	SBLSHC	34
	IF(R(KA).GT.0.)GOTO 35	SBLSHC	35
	T(KA)=0.1	SBLSHC	36
	US(KA)=0	SBLSHC	37
	P(KA)=0	SBLSHC	38
	DP(KA)=0	SBLSHC	39
C	SINGULAR VALUES AT R=0 REPLACED BY ZERO	SBLSHC	40
	GOTO 41	SBLSHC	41
35	T(KA)=TFACT*R(KA)**EX	SBLSHC	42
	US(KA)=R(KA)/(T(KA)*EX)	SBLSHC	43
	P(KA)=2.0*ROC*US(KA)**2/(1.0+GAM)	SBLSHC	44
	UP(KA)=2.0*US(KA)/(1.0+GAM)	SBLSHC	45
	RO(KA)=ROFACT	SBLSHC	46
	DP(KA)=ROFACT*UP(KA)**2*0.5	SBLSHC	47
41	IF(KA.EQ.NR)GOTO 55	SBLSHC	48
	R(KA+1)=R(KA)+(RMAX-R(1))/FLOAT(NR-1)	SBLSHC	49
	IF(R(KA+1).GT.R(KA))GOTO 55	SBLSHC	50
	NBAD=2	SBLSHC	51
	PRINT 45,NBAD,RMAX,RMIN	SBLSHC	52
	RETURN	SBLSHC	53
45	FORMAT(1H0,10X,30HRETURN FROM SBLSHCK WITH NBAD=,I2,	SBLSHC	54
	A14H BECAUSE RMAX=,1PE12.5,20H IS NOT LARGER THAN,	SBLSHC	55
	B6H RMIN=,1PE12.5)	SBLSHC	56
55	CONTINUE	SBLSHC	57



RETURN  
END

SBLSHC58  
SBLSHC59

C	SUBROUTINE SBLPROF(T,RMIN,RMAX,NR,R,P,UP,RO,DP,NBAD)	**** 1
C	THIS COMPUTES THE FLOW PROFILE AT TIME T BETWEEN RMIN AND RMAX	SBLPRO 2
C	T = TIME AFTER THE EXPLOSION	SBLPRO 3
C	RMIN,RMAX = PROFILE LIMITS	SBLPRO 4
C	NR = NUMBER OF NODES TO BE COMPUTED	SBLPRO 5
C	R = DISTANCE FROM THE EXPLOSION	SBLPRO 6
C	P = PRESSURE	SBLPRO 7
C	UP = PARTICLE VELOCITY	SBLPRO 8
C	RO = DENSITY	SBLPRO 9
C	DP = DYNAMIC PRESSURE = 0.5*RO*UP**2	SBLPRO10
C	NBAD = ERROR RETURN INDICATED BY NBAD,NE.0	SBLPRO11
	DIMENSION R(1),P(1),UP(1),RO(1),DP(1)	SBLPRO12
	COMMON/SBLCOM/NC,GAM,ROC,CHEN,V(2),A,B(2),C(5),PC(3),AK,NGOOD	SBLPRO13
	Y(D)=(V(2)/D)**C(1)*((D-V(1))/(V(2)-V(1)))**C(2)	SBLPRO14
A	*((1.-A*D)/(1.-A*V(2)))**C(3)	SBLPRO15
	G(D)=(D/V(2))**C(1)*FLOAT(NC)*((V(1)*GAM-D)/(V(1)*GAM-V(2)))**C(	SBLPRO16
A	4)*((1.-A*D)/(1.-A*V(2)))**C(5)	SBLPRO17
	H(D)=((D-V(1))/(V(2)-V(1)))**C(2)	SBLPRO18
A	*((V(1)*GAM-D)/(V(1)*GAM-V(2)))**C(4)-1.)	SBLPRO19
B	*((1.-A*D)/(1.-A*V(2)))**C(5)-2.*C(3)	SBLPRO20
	U(D)=V(1)+(V(2)-V(1))*D**C(2)	SBLPRO21
	IF(NGOOD,NE.0) GOTO 10	SBLPRO22
	NBAD=11	SBLPRO23
	PRINT 11,NBAD	SBLPRO24
	RETURN	SBLPRO25
11	FORMAT(1H0,10X,30HRETURN FROM SBLPROF WITH NBAD=,I3,	SBLPRO26
	A26H AND WITHOUT COMPUTATIONS.,/,1H,10X,	SBLPRO27
	B34HSUBROUTINE SBLPREP MUST BE CALLED ,	SBLPRO28
	C39H BEFORE OTHER SBL-ROUTINES CAN BE USED.,/)	SBLPRO29
10	NBAD=0	SBLPRO30
	IF(T.GT.0.0)GOTO 25	SBLPRO31
	NBAD=1	SBLPRO32
	PRINT 15,NBAD,T	SBLPRO33
	RETURN	SBLPRO34
15	FORMAT(1H0,10X,30HRETURN FROM SBLPROF WITH NBAD=,I2,	SBLPRO35
	A11H BECAUSE T=,1PE12.5)	SBLPRO36
25	IF(NR.GE.1)GOTO 45	SBLPRO37
	NBAD=2	SBLPRO38
	PRINT 35,NBAD,NR	SBLPRO39
	RETURN	SBLPRO40
35	FORMAT(1H0,10X,30HRETURN FROM SBLPROF WITH NBAD=,I2,	SBLPRO41
	A12H BECAUSE NR=,I4)	SBLPRO42
45	ANT=NC+2	SBLPRO43
	RS=AK*(CHEN/ROC)**(1./ANT)*T**(2./ANT)	SBLPRO44
	USHCK=2.0*RS/(T*ANT)	SBLPRO45
	US=2.0*USHCK/(1.+GAM)	SBLPRO46
	PS=2.0*ROC*USHCK**2/(1.+GAM)	SBLPRO47
	LOS=(GAM+1.0)*ROC/(GAM-1.0)	SBLPRO48
	FNI=AMAX1(RMIN,0.0)	SBLPRO49
	RMI=AMIN1(RMI,RS)	SBLPRO50
	RMA=AMIN1(RMAX,RS)	SBLPRO51
C	NEXT FIND PARAMETER V CORRESPONDING TO RMI	SBLPRO52
	X1=0.0	SBLPRO53
	F1=0.0	SBLPRO54
	X2=1.0	SBLPRO55
	F2=1.0	SBLPRO56
	ICNT=0	SBLPRO57

IF(RMI.GT.0.)GOTO 55	SBLPR058
VMI=0.0	SBLPR059
GOTO 85	SBLPR060
55 IF(RMI.LT.RS)GOTO 65	SBLPR061
VMI=1.0	SBLPR062
GOTO 35	SBLPR063
65 X3=X2-(X2-X1)*(F2-RMI/RS)/(F2-F1)	SBLPR064
XW=W(X3)	SBLPR065
F3=Y(XW)	SBLPR066
IF(ABS(F3-RMI/RS).LT.1.0E-7*RMI/RS) GOTO 75	SBLPR067
ICNT=ICNT+1	SBLPR068
IF(ICNT.GT.20)GOTO 75	SBLPR069
IF(F3-RMI/RS) 56, 75, 58	SBLPR070
56 X1=X3	SBLPR071
F1=F3	SBLPR072
GOTO 65	SBLPR073
58 X2=X3	SBLPR074
F2=F3	SBLPR075
GOTO 65	SBLPR076
75 VMI=X3	SBLPR077
85 R(1)=RS*Y(W(VMI))	SBLPR078
P(1)=PS*G(W(VMI))	SBLPR079
UP(1)=US*(R(1)/RS)*(W(VMI)/V(2))	SBLPR080
RO(1)=ROS*H(W(VMI))	SBLPR081
DP(1)=RO(1)*UP(1)**2*0.5	SBLPR082
IF(NR.EQ.1)RETURN	SBLPR083
IF(RMA.GT.RMI)GOTO 105	SBLPR084
NBAD=3	SBLPR085
PRINT 95,NBAD,RMAX	SBLPR086
RETURN	SBLPR087
95 FORMAT(1H0,10X,30HRETURN FROM SBLPROF WITH NBAD=,I2,	SBLPR088
14H BECAUSE RMAX=,1PE12.5,18H IS OUTSIDE RANGE)	SBLPR089
105 IF(RMA.LT.RS)GOTO 115	SBLPR090
VMA=1.00	SBLPR091
GOTO 145	SBLPR092
C NOW COMPUTE PARAMETER VMA CORRESPONDING TO RMA	SBLPR093
115 ICNT=0	SBLPR094
X1=VMI	SBLPR095
F1=Y(W(X1))	SBLPR096
X2=1.0	SBLPR097
F2=1.0	SBLPR098
125 X3=X2-(X2-X1)*(F2-RMA/RS)/(F2-F1)	SBLPR099
XW=W(X3)	SBLPR100
F3=Y(XW)	SBLPR101
IF(ABS(F3-RMA/RS).LT.1.0E-7 *RMA/RS) GOTO 135	SBLPR102
ICNT=ICNT+1	SBLPR103
IF(ICNT.GT.20)GOTO 135	SBLPR104
IF(F3-RMA/RS)126,135,128	SBLPR105
126 X1=X3	SBLPR106
F1=F3	SBLPR107
GOTO 125	SBLPR108
128 X2=X3	SBLPR109
F2=F3	SBLPR110
GOTO 125	SBLPR111
135 VMA=X3	SBLPR112
145 DO 155 KA=2,NR	SBLPR113
VK=VMI+(VMA-VMI)*FLOAT(KA-1)/FLOAT(NR-1)	SBLPR114

```

VK=W(VK)
F(KA)=RS*Y(VK)
P(KA)=PS*G(VK)
UP(KA)=US*(F(KA)/RS)*(VK/V(2))
RO(KA)=ROS*H(VK)
DP(KA)=RO(KA)*UP(KA)**2*0.5
155 CONTINUE
RETURN
END

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SBLPR115
SBLPR116
SBLPR117
SBLPR118
SBLPR119
SBLPR120
SBLPR121
SBLPR122
SBLPR123

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SUBROUTINE SBLHIST(R,TMIN,TMAX,NR,T,P,UP,RO,DP,NBAD)
C THIS COMPUTES THE FLOW HISTORY AT DISTANCE R AND TIME BETWEEN TMIN,TMSBLHIS 1
C R = DISTANCE SBLHIS 2
C TMIN, TMAX = HISTORY LIMITS SBLHIS 3
C NR = NUMBER OF NODES TO BE COMPUTED SBLHIS 4
C T = TIME SBLHIS 5
C P = PRESSURE SBLHIS 6
C UP = PARTICLE VELOCITY SBLHIS 7
C RO = DENSITY SBLHIS 8
C DP = DYNAMIC PRESSURE = 0.5*RO*UP**2 SBLHIS 9
C NBAD = ERROR RETURN IS INDICATED BY NBAD.NE.0 SBLHIS10
C DIMENSION T(1),P(1),UP(1),RO(1),DP(1) SBLHIS11
C COMMON/SBLCON/NC,GAM,RDC,CHEN,V(2),A,B(2),C(5),PC(3),AK,NGOOD SBLHIS12
Y(D)=(V(2)/D)**C(1)*((D-V(1))/(V(2)-V(1)))**C(2) SBLHIS13
A *((1.-D*A)/(1.-V(2)*A))**C(3) SBLHIS14
C(D)=(D/V(2))**C(1)*FLOAT(NC)*((V(1)*GAM-D)/(V(1)*GAM-V(2)))**C( SBLHIS15
A 4)*((1.-D*A)/(1.-V(2)*A))**C(2)*C(5) SBLHIS16
H(D)=((D-V(1))/(V(2)-V(1)))**C(1)-2.*C(2)) SBLHIS17
A *((V(1)*GAM-D)/(V(1)*GAM-V(2)))**C(4)-1.) SBLHIS18
B *((1.-D*A)/(1.-V(2)*A))**C(2)*C(5)-2.*C(3)) SBLHIS19
W(D)=V(1)+(V(2)-V(1))*D**C(1)/C(2) SBLHIS20
IF(NGOOD.NE.0) GOTO 10 SBLHIS21
NBAD=11 SBLHIS22
PRINT 11,NBAD SBLHIS23
RETURN SBLHIS24
11 FORMAT(1H0,10X,30HRETURN FROM SBLHIST WITH NBAD=,I3, SBLHIS25
A26H AND WITHOUT COMPUTATIONS.,/,1H,10X, SBLHIS26
B34HSUBROUTINE SBLPREP MUST BE CALLED , SBLHIS27
C39H BEFORE OTHER SBL-ROUTINES CAN BE USED.,/) SBLHIS28
10 NBAD=0 SBLHIS29
IF(R.GT.0.)GOTO 25 SBLHIS30
NBAD=1 SBLHIS31
PRINT 15,NBAD,R SBLHIS32
RETURN SBLHIS33
15 FORMAT(1H0,10X,30HRETURN FROM SBLHIST WITH NBAD=,I2, SBLHIS34
A11H BECAUSE R=,1P(12.5) SBLHIS35
25 IF(NR.GE.1)GOTO 27 SBLHIS36
NBAD=2 SBLHIS37
PRINT 26,NBAD,NR SBLHIS38
RETURN SBLHIS39
26 FORMAT(1H0,10X,30HRETURN FROM SBLHIST WITH NBAD=,I2, SBLHIS40
A12H BECAUSE NR=,I4) SBLHIS41
27 ANT=0.5*FLOAT(2+NC) SBLHIS42
ROS=(GAM+1.)*RDC/(GAM-1.) SBLHIS43
TS=SQRT(RDC/CHEN)*(R/AK)**ANT SBLHIS44
C THIS IS SHOCK ARRIVAL TIME AT R SBLHIS45
TNI=ANAL(TMIN,TS) SBLHIS46
IF(TNI.GT.TS)GOTO 35 SBLHIS47
VMI=1.0 SBLHIS48
GOTO 65 SBLHIS49
C NEXT COMPUTE PARAMETER VMI CORRESPONDING TO TNI SBLHIS50
35 ICNT=0 SBLHIS51
DF=(TS/TNI)**(1./ANT) SBLHIS52
X1=0.0 SBLHIS53
F1=0 SBLHIS54
X2 =1.0 SBLHIS55
F2=1. SBLHIS56

```

43	X3=X2-(X2-X1)*(F2-DF)/(F2-F1)	SBLHIS58
	XW=W(X3)	SBLHIS59
	F3=Y(XW)	SBLHIS60
	IF (ABS(F3-DF).LT.1.0E-7 *DF)GOTO 55	SBLHIS61
	ICNT=ICNT+1	SBLHIS62
	IF (ICNT.GT.20)GOTO 55	SBLHIS63
	IF (F3-DF)46,55,48	SBLHIS64
46	X1=X3	SBLHIS65
	F1=F3	SBLHIS66
	GOTO 45	SBLHIS67
48	X2=X3	SBLHIS68
	F2=F3	SBLHIS69
	GOTO 45	SBLHIS70
55	VMI=X3	SBLHIS71
65	XW=W(VMI)	SBLHIS72
	YKA=Y(XW)	SBLHIS73
	T(1)=TS*YKA**(-ANT)	SBLHIS74
	USHCK=(R/YKA)/(T(1)*ANT)	SBLHIS75
	P(1)=2.*(ROC/(1.+GAM))*USHCK**2*G(XW)	SBLHIS76
	UP(1)=R*XW/T(1)	SBLHIS77
	RO(1)=ROS*H(XW)	SBLHIS78
	DP(1)=RO(1)*UP(1)**2*0.5	SBLHIS79
	IF (NR.EQ.1)RETURN	SBLHIS80
	IF (TMAX.GT.THI)GOTO 85	SBLHIS81
	NBAD=3	SBLHIS82
	PRINT 75,NBAD,TMAX	SBLHIS83
	RETURN	SBLHIS84
75	FORMAT(1H0,10X,30HRETURN FROM SBLHIST WITH NBAD=,I4,	SBLHIS85
	A14H BECAUSE TMAX=,1PE12.5,17H IS OUTSIDE RANGE)	SBLHIS86
85	ICNT=0	SBLHIS87
	DF=(TS/TMAX)**(1./ANT)	SBLHIS88
C	NOW FIND VMA CORRESPONDING TO TIME=TMAX	SBLHIS89
	X1=VMI	SBLHIS90
	F1=Y(W(VMI))	SBLHIS91
	X2=0.0	SBLHIS92
	F2=0.0	SBLHIS93
95	X3=X2-(X2-X1)*(F2-DF)/(F2-F1)	SBLHIS94
	XW=W(X3)	SBLHIS95
	F3=Y(XW)	SBLHIS96
	IF (ABS(F3-DF).LT.DF*1.0E-7 )GOTO 105	SBLHIS97
	ICNT=ICNT+1	SBLHIS98
	IF (ICNT.GT.20)GOTO 105	SBLHIS99
	IF (F3-DF)96,105,98	SBLHI100
96	X2=X3	SBLHI101
	F2=F3	SBLHI102
	GOTO 95	SBLHI103
98	X1=X3	SBLHI104
	F1=F3	SBLHI105
	GOTO 95	SBLHI106
105	VMA=X3	SBLHI107
	DO 115 KA=2,NR	SBLHI108
	VK=VMI+(VMA-VMI)*FLOAT(KA-1)/FLOAT(NR-1)	SBLHI109
	XW=W(VK)	SBLHI110
	YKA=Y(XW)	SBLHI111
	T(KA)=TS*YKA**(-ANT)	SBLHI112
	RS=R/YKA	SBLHI113
	USHCK=RS/(T(KA)*ANT)	SBLHI114

P(KA)=2.\*(RDC/(1.+GAM))\*USHCK\*\*2\*G(XH)  
UP(KA)=R\*XW/T(KA)  
RO(KA)=ROS\*H(XW)  
DP(KA)=RO(KA)\*UP(KA)\*\*2\*0.5  
115 CONTINUE  
RETURN  
END

SBLHI115  
SBLHI116  
SBLHI117  
SBLHI118  
SBLHI119  
SBLHI120  
SBLHI121



C	SUBROUTINE SBLPATH(RZ,RMAX,TMAX,NR,R,T,P,UP,RO,DP,NBAD)	**** 1
C	THIS COMPUTES A PARTICLE PATH STARTING ON THE SHOCK AT R=RZ	SBLPAT 2
C	RZ = DISTANCE OF FIRST NODE FROM EXPLOSION, LOCATED ON THE SHOCK	SBLPAT 3
C	RMAX,TMAX = PATH ENDS WHEN EITHER OF THESE IS REACHED	SBLPAT 4
C	NR = NUMBER OF NODES TO BE COMPUTED	SBLPAT 5
C	R = DISTANCE FROM THE EXPLOSION	SBLPAT 6
C	T = TIME AFTER EXPLOSION	SBLPAT 7
C	P = PRESSURE	SBLPAT 8
C	UP = PARTICLE VELOCITY	SBLPAT 9
C	RO = DENSITY	SBLPAT10
C	DP = DYNAMIC PRESSURE = 0.5*RO*UP**2	SBLPAT11
C	NBAD = ERROR RETURN WILL BE INDICATED BY NBAD.NE.0	SBLPAT12
	DIMENSION R(1),T(1),P(1),UP(1),RO(1),DP(1)	SBLPAT13
	COMMON/SBLCON/NC,GAM,ROC,CHEN,V(2),A,B(2),C(5),PC(3),AK,NGOOD	SBLPAT14
	Y(D)=(V(2)/D)**C(1)*((D-V(1))/(V(2)-V(1)))**C(2)	SBLPAT15
	A *((1.-D*A)/(1.-V(2)*A))**C(3)	SBLPAT16
	G(D)=(D/V(2))*((C(1)*FLOAT(NC))*((V(1)*GAM-D)/(V(1)*GAM-V(2))))**C(	SBLPAT17
	4)*((1.-D*A)/(1.-V(2)*A))**C(5))	SBLPAT18
	H(D)=((D-V(1))/(V(2)-V(1)))**C(2)	SBLPAT19
	A *((V(1)*GAM-D)/(V(1)*GAM-V(2)))**C(4)-1.)	SBLPAT20
	B *((1.-D*A)/(1.-V(2)*A))**C(5)-2.*C(3))	SBLPAT21
	W(D)=(D/V(2))*((V(1)*GAM-D)/(V(1)*GAM-V(2)))**PC(1)	SBLPAT22
	A *((D-V(1))/(V(2)-V(1)))**PC(2)*((1.-D*A)/(1.-V(2)*A))**PC(3)	SBLPAT23
	RSF(D)=AK*(CHEN/ROC)**(1./FLOAT(2+NC))*D**(2./FLOAT(2+NC))	SBLPAT24
	X(D)=V(1)+(V(2)-V(1))*D**(-1./PC(2))	SBLPAT25
	IF(NGOOD.NE.0) GOTO 10	SBLPAT26
	NBAD=11	SBLPAT27
	PRINT 11,NBAD	SBLPAT28
	RETURN	SBLPAT29
11	FORMAT(1H0,10X,30HRETURN FROM SBLPATH WITH NBAD=,I3,	SBLPAT30
	A26H AND WITHOUT COMPUTATIONS.,/,1H,10X,	SBLPAT31
	B34HSUBROUTINE SBLPREP MUST BE CALLED ,	SBLPAT32
	C39H BEFORE OTHER SBL-ROUTINES CAN BE USED.,/)	SBLPAT33
10	NBAD=0	SBLPAT34
	IF(RZ.GT.0.)GOTO 25	SBLPAT35
	NBAD=1	SBLPAT36
	PRINT 15,NBAD,RZ	SBLPAT37
	RETURN	SBLPAT38
15	FORMAT(1H0,10X,30HRETURN FROM SBLPATH WITH NBAD=,I2,	SBLPAT39
	A12H BECAUSE RZ=,1PE12.5)	SBLPAT40
25	IF(NR.GT.0)GOTO 45	SBLPAT41
	NBAD=2	SBLPAT42
	PRINT 35,NBAD,NR	SBLPAT43
	RETURN	SBLPAT44
35	FORMAT(1H0,10X,30HRETURN FROM SBLPATH WITH NBAD=,I2,	SBLPAT45
	A12H BECAUSE NR=,I4)	SBLPAT46
45	R(1)=RZ	SBLPAT47
	ANT=2+NC	SBLPAT48
	T(1)=AK**(-0.5*ANT)*SQRT(ROC/CHEN)*RZ**(0.5*ANT)	SBLPAT49
	USHCK=(2.0/ANT)*RZ/T(1)	SBLPAT50
	P(1)=(2.0/(GAM+1.))*ROC*USHCK**2	SBLPAT51
	UP(1)=(2.0/(GAM+1.))*USHCK	SBLPAT52
	RO(1)=((GAM+1.)/(GAM-1.))*ROC	SBLPAT53
	DP(1)=RO(1)*UP(1)**2*0.5	SBLPAT54
	IF(NR.EQ.1)RETURN	SBLPAT55
	IF(TMAX.GT.T(1).AND.RMAX.GT.R(1))GOTO 65	SBLPAT56
	NBAD=3	SBLPAT57

```

      PRINT 55,NBAD,TMAX,T(1),RMAX,R(1)
      RETURN
55  FORMAT(1H0,10X,30HRETURN FROM SBLPATH WITH NBAD=,I2,
      A14H BECAUSE TMAX=,1PE12.5,13H .LE. TSHOCK=,1PE12.5,/,
      B1H ,10X,16HOR BECAUSE RMAX=,1PE12.5,9H .LE. RZ=,1PE12.5)
65  ICNT=0
C   NOW COMPUTE PARAMETER VMA CORRESPONDING TO TMAX
      DT=T(1)/TMAX
      X1=0.0
      F1=0.0
      X2=1.0
      F2=1.0
75  X3=X2-(X2-X1)*(F2-DT)/(F2-F1)
      F3=1./U(X(X3))
      IF(ABS(F3-DT).LT.1.0E-7*DT) GOTO 85
      ICNT=ICNT+1
      IF(ICNT.GT.20)GOTO 85
      IF(F3-DT) 76,85,78
76  X1=X3
      F1=F3
      GOTO 75
78  X2=X3
      F2=F3
      GOTO 75
85  VTHAX=X3
      RTMAX=RSF(TMAX)*Y(X(VTHAX))
      IF(RTMAX.LT.RMAX)GOTO 105
C   BRANCH IF END OF PATH DETERMINED BY TMAX
C   ELSE COMPUTE PARAMETER CORRESPONDING TO RMAX
      ICNT=0
      X1=VTHAX
      F1=RTMAX
      X2=1.0
      F2=RZ
95  X3=X2-(X2-X1)*(F2-RMAX)/(F2-F1)
      DT=T(1)*W(X(X3))
      F3=RSF(DT)*Y(X(X3))
      IF(ABS(F3-RMAX).LT.1.0E-7*RMAX) GOTO 105
      ICNT=ICNT+1
      IF(ICNT.GT.20)GOTO 105
      IF(F3-RMAX) 96,105,98
96  X2=X3
      F2=F3
      GOTO 95
98  X1=X3
      F1=F3
      GOTO 95
105 YMAX=1./X3
      YMIN=1.0
      DO 135 KA=2,NR
      YK=YMIN+(YMAX-YMIN)*FLOAT(KA-1)/FLOAT(NR-1)
      VK=V(1)+(V(2)-V(1))*YK**(1./PC(2))
      T(KA)=T(1)*U(VK)
      RS=RSF(T(KA))
      R(KA)=RS*Y(VK)
      USHCK=(RS/T(KA))*2.0/FLOAT(2+NC)
      P(KA)=(2.0/(1.+GAM))*ROC*USHCK**2*G(VK)

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SBLPAT58
SBLPAT59
SBLPAT60
SBLPAT61
SBLPAT62
SBLPAT63
SBLPAT64
SBLPAT65
SBLPAT66
SBLPAT67
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SBLPAT95
SBLPAT96
SBLPAT97
SBLPAT98
SBLPAT99
SBLPA100
SBLPA101
SBLPA102
SBLPA103
SBLPA104
SBLPA105
SBLPA106
SBLPA107
SBLPA108
SBLPA109
SBLPA110
SBLPA111
SBLPA112
SBLPA113
SBLPA114

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UP(KA)=R(KA)*VK/T(KA)
RO(KA)=((GAM+1.)/(GAM-1.))*ROD*H(VK)
DP(KA)=RO(KA)*UP(KA)**2*0.5
135 CONTINUE
RETURN
END
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SBLPA115
SBLPA116
SBLPA117
SBLPA118
SBLPA119
SBLPA120
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SUBROUTINE SBLMACH(RZ,TZ,RMIN,RMAX,TMIN,TMAX,NR,
A R,T,P,UP,RO,DP,HEAD)
C THIS COMPUTES MACH-LINES AND CORRESPONDING FLOW FIELD
C RZ,TZ = INITIAL POINT OF THE MACH LINES. EITHER RZ OR TZ MUST
C BE ZERO. IF RZ=0, THEN COMPUTATION WILL START
C AT THE EXPLOSION AND TIME TZ.
C IF TZ=0, THEN COMPUTATION WILL START ON THE
C SHOCK AT DISTANCE RZ AND CORRESPONDING TIME
C RMIN,---,TMAX = LIMITS FOR THE COMPUTATIONS
C NR(2) = NUMBER OF NODES TO BE COMPUTED.
C NR(1) CORRESPONDS TO CHARACTERISTIC WITH DT/DR.GT.0
C R(2,1) = DISTANCE FROM THE EXPLOSION
C T(2,1) = TIME FROM THE EXPLOSION
C P(2,1) = PRESSURE
C UP(2,1) = PARTICLE VELOCITY
C RO(2,1) = DENSITY
C DP(2,1) = DYNAMIC PRESSURE = 0.5*P0*UP**2
C NBAD = ERROR RETURN WILL BE INDICATED BY NBAD.NE.0
C DIMENSION NR(2),R(2,1),T(2,1),P(2,1),RO(2,1),DP(2,1),UP(2,1)
C COMMON/SBLCON/NC,GAM,ROC,CHEN,V(2),A,B(2),C(5),PC(3),AK,NGOOD
C Y(D)=(V(2)/D)**C(1)**((D-V(1))/(V(2)-V(1)))**C(2)
C A *((1.-D*A)/(1.-V(2)*A))**C(3)
C G(D)=(D/V(2))**((C(1)*FLOAT(NC))**((V(1)*GAM-D)/(V(1)*GAM-V(2)))**C(
C 4)*((1.-D*A)/(1.-V(2)*A))**C(5))
C H(D)=((D-V(1))/(V(2)-V(1)))**C(2)**(1.-2.*C(2))
C A *((V(1)*GAM-D)/(V(1)*GAM-V(2)))**C(5)-2.*C(3))
C RSF(D)=AK*((CHEN/ROC)*D**2)**(1./FLOAT(2+NC))
C Z1(D)=A*D/(1.-A*D)
C ZA(D)=(B(1)+B(2)*D*A)/((1.-D*A)*V(1)*A)
C ZB(D)=AMAX1(-1.0,AMIN1(1.0,ZA(D)))
C Z2(D)=EXP(ASIN(ZB(D))/SQRT(B(1)+B(2)))
C XF(D)=V(1)+(V(2)-V(1))*D**C(2)
C IF(NGOOD.NE.0) GOTO 10
C NBAD=11
C PRINT 11,NBAD
C RETURN
11 FORMAT(1H0,10X,30HRETURN FROM SBLMACH WITH NBAD=,I3,
A26H AND WITHOUT COMPUTATIONS.,/ ,1H ,10X,
B34HSUBROUTINE SBLPREP MUST BE CALLED ,
C39H BEFORE OTHER SBL-ROUTINES CAN BE USED.,/)
10 NBAD=0
C Z2V1=EXP(-1.5707963268/SQRT(B(1)+B(2)))
C IF(NR(1).GE.1.OR.NR(2).GE.1)GOTO 25
C NBAD=1
12 PRINT 15,NBAD
C RETURN
15 FORMAT(1H0,10X,30HRETURN FROM SBLMACH WITH NBAD=,I2,
A39H BECAUSE ARGUMENTS ARE NOT WITHIN RANGE)
25 IF(RZ.EQ.0..AND.TZ.GT.0.)GOTO 35
C IF(RZ.GT.0..AND.TZ.EQ.0.)GOTO 45
C NBAD=2
C GOTO 12
C IN THIS CASE THE INITIAL NODE IS AT THE EXPLOSION
35 TRAT=(Z1(V(1))*Z2V1)/(Z1(V(2))*Z2(V(2)))
C NUP=1
C IF(TRAT.GT.1.)NUP=-1
C ANUP=TZ/(Z1(V(1))*Z2V1**NUP)

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TDUP=TZ	SBLMAC58
TFUP=AHUP*Z1(V(2))*Z2(V(2))*NUP	SBLMAC59
RFUP=RSF(TFUP)	SBLMAC60
AMDOWN=TZ/(Z1(V(1))*Z2V1**(-NUP))	SBLMAC61
TBDOWN=TZ	SBLMAC62
TFDOWN=AMDOWN*Z1(V(2))*Z2(V(2))*(-NUP)	SBLMAC63
RFDOWN=RSF(TFDOWN)	SBLMAC64
C THIS ESTABLISHED END POINTS OF BOTH CHARACTERISTICS.	SBLMAC65
C NOW GO TO FIND OUT WHICH SEGMENT SHOULD BE COMPUTED	SBLMAC66
GO TO 55	SBLMAC67
45 TS=SQRT(ROC/CHEN)*(RZ/AK)**(0.5*FLOAT(2+NC))	SBLMAC68
C IN THIS CASE A POINT ON THE SHOCK IS SPECIFIED	SBLMAC69
TRAT=(Z1(V(1))*Z2V1)/(Z1(V(2))*Z2(V(2)))	SBLMAC70
NUP=1	SBLMAC71
IF(TRAT.GT.1.0)NUP=-1	SBLMAC72
ANUP=TS/(Z1(V(2))*Z2(V(2))*NUP)	SBLMAC73
TBUP=AHUP*Z1(V(1))*Z2V1**NUP	SBLMAC74
TFUP=TS	SBLMAC75
RFUP=RZ	SBLMAC76
AMDOWN=TS/(Z1(V(2))*Z2(V(2))*(-NUP))	SBLMAC77
TBDOWN=AMDOWN*Z1(V(1))*Z2V1**(-NUP)	SBLMAC78
TFDOWN=TS	SBLMAC79
RFDOWN=RZ	SBLMAC80
C AT 55 START WORK ON UPWARD CHARACTERISTIC	SBLMAC81
55 IF(NR(1).LE.0)GO TO 205	SBLMAC82
C BRANCH TO DOWNWARD CHARACTERISTIC	SBLMAC83
IF(RFUP.LT.RMIN.OR.TFUP.LT.THIN)GO TO 65	SBLMAC84
IF(0.6.GT.RMAX.OR.TBUP.GT.TMAX)GO TO 65	SBLMAC85
GO TO 75	SBLMAC86
65 NR(1)=0	SBLMAC87
GO TO 205	SBLMAC88
C NOW FIND INTERSECTION WITH RMIN	SBLMAC89
75 X1=0.1	SBLMAC90
F1=0.1	SBLMAC91
X2=1.0	SBLMAC92
T2=TFUP	SBLMAC93
F2=RFUP	SBLMAC94
ICNT=0	SBLMAC95
RM=AMAX1(RMIN,0.)	SBLMAC96
85 X3=X2-(X2-X1)*(F2-RM)/(F2-F1)	SBLMAC97
X3=AMAX1(0.,X3)	SBLMAC98
IF(X3.LE.0.)T3=AHUP*Z1(XF(X3))*Z2V1**NUP	SBLMAC99
IF(X3.GT.0.)T3=AMUP*Z1(XF(X3))*Z2(XF(X3))*NUP	SBLMA100
F3=RSF(T3)*Y(XF(X3))	SBLMA101
IF(ABS(F3-RM).LT.1.0E-7*RFDOWN) GO TO 95	SBLMA102
ICNT=ICNT+1	SBLMA103
IF(ICNT.GT.20)GO TO 95	SBLMA104
IF(F3-RM) 86,95,88	SBLMA105
86 X1=X3	SBLMA106
F1=F3	SBLMA107
T1=T3	SBLMA108
GO TO 85	SBLMA109
88 X2=X3	SBLMA110
F2=F3	SBLMA111
T2=T3	SBLMA112
GO TO 85	SBLMA113
95 X3=XF(X3)	SBLMA114



IF(T3-TMAX) 115,105,65	SBLMA115
105 NR(1)=1	SBLMA116
GOTO 135	SBLMA117
115 IF(T3.GE.TMIN) GOTO 135	SBLMA118
C BRANCH IF FIRST NODE WAS FOUND. ELSE GET INTERSECTION WITH TMIN	SBLMA119
X1=X3	SBLMA120
F1=T3	SBLMA121
X2=V(2)	SBLMA122
F2=TFUP	SBLMA123
T2=TFUP	SBLMA124
ICNT=0	SBLMA125
125 X3=X2-(X2-X1)*(F2-TMIN)/(F2-F1)	SBLMA126
T3=ANUP*Z1(X3)*Z2(X3)**NUP	SBLMA127
F3=T3	SBLMA128
IF(ABS(F3-TMIN).LT.1.0E-7*TBUP) GOTO 135	SBLMA129
ICNT=ICNT+1	SBLMA130
IF(ICNT.GT.20)GOTO 135	SBLMA131
IF(F3-TMIN) 126,135,128	SBLMA132
126 X1=X3	SBLMA133
T1=T3	SBLMA134
F1=F3	SBLMA135
GOTO 125	SBLMA136
128 X2=X3	SBLMA137
T2=T3	SBLMA138
F2=F3	SBLMA139
GOTO 125	SBLMA140
135 VIN=X3	SBLMA141
RS=RSF(T3)	SBLMA142
T(1,1)=T3	SBLMA143
R(1,1)=RS*Y(VIN)	SBLMA144
IF(R(1,1).GT.RMAX)GOTO 65	SBLMA145
USHCK=(2./FLOAT(2+NC))*RS/T(1,1)	SBLMA146
UP(1,1)=(R(1,1)/T(1,1))*VIN	SBLMA147
P(1,1)=(2./((GAM+1.))*RDC*USHCK**2*G(VIN)	SBLMA148
RD(1,1)=((GAM+1.)/(GAM-1.))*RDC*H(VIN)	SBLMA149
DP(1,1)=RD(1,1)*UP(1,1)**2*0.5	SBLMA150
IF(NR(1).EQ.1)GOTO 205	SBLMA151
C BRANCH TO COMPUTATION OF DOWNWARD CHARACTERISTIC	SBLMA152
C NOW FIND END POINT OF CURVE	SBLMA153
IF(R(1,1).EQ.RMAX)GOTO 177	SBLMA154
X1=((VIN-V(1))/(V(2)-V(1)))*C(2)	SBLMA155
F1=R(1,1)	SBLMA156
X2=1.0	SBLMA157
T2=TFUP	SBLMA158
F2=RFUP	SBLMA159
ICNT=0	SBLMA160
RM=AMIN1(RMAX,RFUP)	SBLMA161
145 X3=X2-(X2-X1)*(F2-RM)/(F2-F1)	SBLMA162
T3=ANUP*Z1(XF(X3))*Z2(XF(X3))**NUP	SBLMA163
F3=RSF(T3)*Y(XF(X3))	SBLMA164
IF(ABS(F3-RM).LT.1.0E-7*RFDOWN) GOTO 155	SBLMA165
ICNT=ICNT+1	SBLMA166
IF(ICNT.GT.20)GOTO 155	SBLMA167
IF(F3-RM) 146,155,148	SBLMA168
146 X1=X3	SBLMA169
T1=T3	SBLMA170
F1=F3	SBLMA171

GOTO 145	SBLMA172
148 X2=X3	SBLMA173
T2=T3	SBLMA174
F2=F3	SBLMA175
GOTO 145	SBLMA176
155 X3=XF(X3)	SBLMA177
IF(T3.LE.TMAX) GOTO 175	SBLMA178
C BRANCH IF END OF CURVE FOUND. ELSE FIND INTERSECTION WITH TMAX	SBLMA179
X1=VIN	SBLMA180
F1=T(1,1)	SBLMA181
X2=X3	SBLMA182
F2=T3	SBLMA183
ICNT=0	SBLMA184
165 X3=X2-(X2-X1)*(F2-TMAX)/(F2-F1)	SBLMA185
T3=AMUP*Z1(X3)*Z2(X3)**NUP	SBLMA186
F3=T3	SBLMA187
IF(ABS(F3-TMAX).LT.1.0E-7 *TMAX)GOTO 175	SBLMA188
ICNT=ICNT+1	SBLMA189
IF(ICNT.GT.20)GOTO 175	SBLMA190
IF(F3-TMAX) 166,175,168	SBLMA191
166 X1=X3	SBLMA192
T1=T3	SBLMA193
F1=F3	SBLMA194
GOTO 165	SBLMA195
168 X2=X3	SBLMA196
T2=T3	SBLMA197
F2=F3	SBLMA198
GOTO 165	SBLMA199
175 VEN=X3	SBLMA200
IF(VEN.GT.VIN)GOTO 185	SBLMA201
177 NR(1)=1	SBLMA202
GOTO 205	SBLMA203
185 KUP=NR(1)	SBLMA204
XIN=((VIN-V(1))/(V(2)-V(1)))*C(2)	SBLMA205
XEN=((VEN-V(1))/(V(2)-V(1)))*C(2)	SBLMA206
DO 195 KA=2,KUP	SBLMA207
X=XIN+(XEN-XIN)*FLOAT(KA-1)/FLOAT(KUP-1)	SBLMA208
X=XF(X)	SBLMA209
X=AMAX1(X,V(1))	SBLMA210
IF(X.LE.V(1)) T(1,KA)=AMUP*Z1(V(1))*Z2V1**NUP	SBLMA211
IF(X.GT.V(1)) T(1,KA)=AMUP*Z1(X)*Z2(X)**NUP	SBLMA212
RS=RSF(T(1,KA))	SBLMA213
R(1,KA)=RS*Y(X)	SBLMA214
USHCK=(2./FLOAT(2+HC))*RS/T(1,KA)	SBLMA215
P(1,KA)=(2./(GAM+1.))*RDC*USHCK**2*G(X)	SBLMA216
UP(1,KA)=X*R(1,KA)/T(1,KA)	SBLMA217
RD(1,KA)=((GAM+1.)/(GAM-1.))*RDC*H(X)	SBLMA218
DP(1,KA)=RD(1,KA)*UP(1,KA)**2*0.5	SBLMA219
195 CONTINUE	SBLMA220
C AT 205 START COMPUTATION OF DOWNWARD CURVE	SBLMA221
205 IF(NR(2).LE.0)RETURN	SBLMA222
IF(RFDOWN.LT.RMIN.OR.TFDOWN.GT.TMAX)GOTO 215	SBLMA223
IF(0.LT.RMAX.OR.TBDOWN.LT.TMIN)GOTO 215	SBLMA224
GOT 225	SBLMA225
215 NR(2)=0	SBLMA226
RETURN	SBLMA227
C AT 225 FIND DOWNWARD INTERSECTION WITH RMIN	SBLMA228



225	X1=0.0	SBLMA229
	F1=0	SBLMA230
	X2=1.0	SBLMA231
	T2=TFDOWN	SBLMA232
	F2=RFDOWN	SBLMA233
	ICNT=0	SBLMA234
	FM=AMAX1(RMIN,0.)	SBLMA235
235	X3=X2-(X2-X1)*(F2-FM)/(F2-F1)	SBLMA236
	X3=AMAX1(0.,X3)	SBLMA237
	IF(X3.LE.0.)T3=AMDOWN*Z1(XF(X3))/Z2V1**NUP	SBLMA238
	IF(X3.GT.0.)T3=AMDOWN*Z1(XF(X3))/Z2(XF(X3))**NUP	SBLMA239
	F3=RSF(T3)*Y(XF(X3))	SBLMA240
	IF(ABS(F3-FM).LT.1.0E-7 *RFDOWN)GOTO 245	SBLMA241
	ICNT=ICNT+1	SBLMA242
	IF(ICNT.GT.20)GOTO 245	SBLMA243
	IF(F3-FM) 236,245,238	SBLMA244
236	X1=X3	SBLMA245
	T1=T3	SBLMA246
	F1=F3	SBLMA247
	GOTO 235	SBLMA248
238	X2=X3	SBLMA249
	T2=T3	SBLMA250
	F2=F3	SBLMA251
	GOTO 235	SBLMA252
245	X3=XF(X3)	SBLMA253
	IF(T3-TMIN)215,255,265	SBLMA254
255	NR(2)=1	SBLMA255
	GOTO 265	SBLMA256
265	IF(T3.LE.TMAX)GOTO 285	SBLMA257
C	BRANCH IF FIRST NODE FOUND	SBLMA258
	X1=X3	SBLMA259
	F1=T3	SBLMA260
	X2=V(2)	SBLMA261
	T2=TFDOWN	SBLMA262
	F2=T2	SBLMA263
	ICNT=0	SBLMA264
275	X3=X2-(X2-X1)*(F2-TMAX)/(F2-F1)	SBLMA265
	T3=AMDOWN*Z1(X3)/Z2(X3)**NUP	SBLMA266
	F3=T3	SBLMA267
	IF(ABS(F3-TMAX).LT.1.0E-7 *TFDOWN)GOTO 285	SBLMA268
	ICNT=ICNT+1	SBLMA269
	IF(ICNT.GT.20)GOTO 285	SBLMA270
	IF(F3-TMAX) 276,285,278	SBLMA271
276	X2=X3	SBLMA272
	T2=T3	SBLMA273
	F2=F3	SBLMA274
	GOTO 275	SBLMA275
278	X1=X3	SBLMA276
	T1=T3	SBLMA277
	F1=F3	SBLMA278
	GOTO 275	SBLMA279
285	VIN=X3	SBLMA280
	RS=RSF(T3)	SBLMA281
	T(2,1)=T3	SBLMA282
	R(2,1)=RS*Y(VIN)	SBLMA283
	IF(R(2,1).GT.RMAX)GOTO 215	SBLMA284
	USHCY=(2./FLOAT(2+NC))*RS/T(2,1)	SBLMA285

P(2,1)=(2.0/(GAM+1.0))*ROC*USHCK**2*G(VIN)	SBLMA286
UP(2,1)=VIN*R(2,1)/T(2,1)	SBLMA287
RO(2,1)=((GAM+1.0)/(GAM-1.0))*ROC*H(VIN)	SBLMA288
DP(2,1)=RO(2,1)*UP(2,1)**2*0.15	SBLMA289
IF(R(2,1).EQ.RMAX)NR(2)=1	SBLMA290
IF(NR(2).EQ.1)RETURN	SBLMA291
C NOW FIND END POINT OF DOWNWARD CURVE	SBLMA292
C FIRST FIND INTERSECTION WITH RMAX	SBLMA293
X1=((VIN-V(1))/(V(2)-V(1)))*C(2)	SBLMA294
F1=R(2,1)	SBLMA295
X2=1.0	SBLMA296
T2=TFDOWN	SBLMA297
F2=RFDOWN	SBLMA298
ICNT=0	SBLMA299
RH=AMIN1(RFDOWN,RMAX)	SBLMA300
295 X3=X2-(X2-X1)*(F2-RH)/(F2-F1)	SBLMA301
T3=ANDOWN*Z1(XF(X3))/Z2(XF(X3))*NUP	SBLMA302
F3=RSF(T3)*Y(XF(X3))	SBLMA303
IF(ABS(F3-RH).LT.1.0E-7 *RFDOWN)GOTO 305	SBLMA304
ICNT=ICNT+1	SBLMA305
IF(ICNT.GT.20)GOTO 305	SBLMA306
IF(F3-RH) 296,305,298	SBLMA307
296 X1=X3	SBLMA308
T1=T3	SBLMA309
F1=F3	SBLMA310
GOTO 295	SBLMA311
298 X2=X3	SBLMA312
T2=T3	SBLMA313
F2=F3	SBLMA314
GOTO 295	SBLMA315
305 X3=XF(X3)	SBLMA316
IF(T3.GE.TMIN)GOTO 325	SBLMA317
C BRANCH IF END POINT FOUND. ELSE FIND INTERSECTION WITH TMIN	SBLMA318
X1=VIN	SBLMA319
F1=T(2,1)	SBLMA320
F2=T3	SBLMA321
X2=X3	SBLMA322
ICNT=0	SBLMA323
315 X3=X2-(X2-X1)*(F2-TMIN)/(F2-F1)	SBLMA324
T3=ANDOWN*Z1(X3)/Z2(X3))*NUP	SBLMA325
F3=T3	SBLMA326
IF(ABS(F3-TMIN).LT.1.0E-7 *TFDOWN)GOTO 325	SBLMA327
ICNT=ICNT+1	SBLMA328
IF(ICNT.GT.20)GOTO 325	SBLMA329
IF(F3-TMIN) 316,325,318	SBLMA330
316 X2=X3	SBLMA331
T2=T3	SBLMA332
F2=F3	SBLMA333
GOTO 315	SBLMA334
318 X1=X3	SBLMA335
T1=T3	SBLMA336
F1=F3	SBLMA337
GOTO 315	SBLMA338
325 VEN=X3	SBLMA339
IF(VEN.GT.VIN)GOTO 335	SBLMA340
NR(2)=1	SBLMA341
RETURN	SBLMA342

C LOOP TO COMPUTE DOWNWARD CHARACTERISTIC

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335 KUP=NR(2)
XIN=((VIN-V(1))/(V(2)-V(1)))*C(2)
XEN=((VEN-V(1))/(V(2)-V(1)))*C(2)
DO 345 KA=2,KUP
X=XIN+(XEN-XIN)*FLOAT(KA-1)/FLOAT(KUP-1)
X=XF(X)
X=AMAX1(X,V(1))
IF(X.LE.V(1)) T(2,KA)=AMDOWN*Z1(V(1))/Z2V1**NUP
IF(X.GT.V(1)) T(2,KA)=AMDOWN*Z1(X)/Z2(X)**NUP
RS=RSF(T(2,KA))
R(2,KA)=RS*Y(X)
USHCK=(2./FLOAT(2+HC))*RS/T(2,KA)
P(2,KA)=(2./GAM+1.0)*ROC*USHCK**2*G(X)
UP(2,KA)=X*R(2,KA)/T(2,KA)
RO(2,KA)=((GAM+1.)/(GAM-1.))*ROC*H(X)
DP(2,KA)=RO(2,KA)*UP(2,KA)**2*0.5
345 CONTINUE
RETURN
END

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